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### **DL** Software and **Lattice Boltzmann Equation (LBE)**

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### The mesoscale

'If life is going to exist in a Universe of this size, then the one thing it cannot afford to have is a sense of proportion.'Douglas Adams, *The Restaurant at the End of the Universe* 

### Getting a sense of (middle) scale

### Mesoscale

- Bridges a 'hinterland' between atomistic and continuum scales
  - From 10nm+, 1ns+
  - Upper limit depends on computing capability
- Many interesting scientific applications operate at mesoscale
  - Approaching engineering (process) scales



- Correct thermodynamics <u>and</u> hydrodynamics needed
  - Parameterisation may involve bottom-up (microscopic) <u>and</u> top-down (macroscopic)
- Models often involve bigger-than-atom particles ('beads')



### Mesoscale modelling approaches

### **Bottom-up: from the microscale**

- Coarse-grain atoms/molecules into beads as 'sub-thermodynamic populations'
- Obtain effective interactions between beads to represent those between atoms

### **Top-down: from the macroscale**

- Reverse-engineer simulation rules to give required behaviours
- Beads are 'carriers of momentum': mesoscopic representation of continuum fluids







### Lattice Boltzmann Equation (LBE)

'He was constantly reminded of how startlingly different a place the world was when viewed from a point only three feet to the left.'

- Douglas Adams, The Salmon of Doubt

### Lattice Boltzmann Equation (LBE)

### Moving from particle dynamics to statistical mechanics

- Still want to consider how particles move around, but at <u>very</u> large scales (cf. atomistic MD) the <u>motion of an individual particle</u> matters very little!
- An approach that can work at mesoscale: using probabilities of finding particles and evolving these over time
  - **Distribution functions**: probabilities of finding particles at given time and space with specific momenta (moments give fluid density, momentum etc.)
  - Distribution functions evolve via **Boltzmann equation**:

Propagation (free-streaming) of particles in system due to velocity field and applied forces

$$\frac{\partial f}{\partial t} + \frac{\mathbf{p}}{m} \cdot \nabla f + \mathbf{F} \cdot \frac{\partial f}{\partial \mathbf{p}} = \left(\frac{\partial f}{\partial t}\right)_{coll}$$

Collisions of particles when encountering each other: governed by e.g. fluid viscosity

- Collisions can be complex, especially if large number of particles involved
- Keeping track of particles and solving Boltzmann equation more difficult if there are too many possible momenta and positions

### Lattice gas cellular automata Antecedent to LBE (and DPD)

- Gas particles stream ('propagate') on a regular lattice along links with set velocities e<sub>i</sub>
- Particles collide as they encounter each other at lattice sites: maintain mass/total momentum, related to viscosity



$$n_i(\mathbf{x} + \mathbf{e}_i \Delta t, t + \Delta t) - n_i(\mathbf{x}, t) = \Delta_i(n_0, n_1, \dots n_N)$$



- Boolean/integer calculations: unconditional numerical stability
- Easy to parallelise
- Multi-physics possible
- Can show connection to Navier-Stokes fluid flow equations

#### Cons

- Statistically noisy: limited viscosities available
- Ensemble-averaging required to see results
- No (automatic) Galilean invariance, anisotropic flows
- Large number of possible collision states
- Only small number of lattice schemes will work



### **LBE: LGCA with statistical mechanics**

### Getting from LGCA to LBE in three easy(ish!) steps

- 1. Substitute Boolean particle occupation numbers  $n_i$  with density distribution functions  $f_i^{[1]}$ 
  - Ensemble-averaging LGCA ab initio
  - Moments of  $f_i$  give macroscopic fluid density, momentum etc.

$$f_i = \langle n_i \rangle$$
  

$$\rho(\mathbf{x}) = \sum_i f_i(\mathbf{x})$$
  

$$\rho(\mathbf{x}) \mathbf{u}(\mathbf{x}) = \sum_i \mathbf{e}_i f_i(\mathbf{x})$$

Eliminates statistical noise of LGCA and need to ensemble-average results to actually see flows(!), widens available viscosity range



1. McNamara and Zanetti, *PRL* **61** (20), 2332–2335 (1988)

### LBE: LGCA with statistical mechanics

### Getting to LBE in three easy(ish!) steps

- 2. Transform collision matrix  $\Delta_i$  into linearised form based on deviation from local equilibrium  $f_i^{eq}$ <sup>[1]</sup>
  - Matrix  $\Omega_{ij}$  considerably smaller than before (e.g. from 2<sup>N</sup> elements to N<sup>2</sup> for N links per grid point)
  - Can now be considered as a **collision operator**: tries to shift each point towards an equilibrium state

$$\begin{split} \Delta_i(f_0, f_1, \dots f_N) &\to \sum_j \Omega_{ij} \left( f_j - f_j^{eq} \right) \\ f_i^{eq} &= f(\rho, \mathbf{u}) \end{split}$$

Big improvement in computational speed, Galilean invariance now guaranteed



1. Succi *et al.*, *Physica D* **47** (1–2), 219–230 (1991)

### LBE: LGCA with statistical mechanics

### Getting to LBE in three easy(ish!) steps

- 3. Simplify collisions even further by setting  $\Omega_{ii}$  to unit diagonal with a relaxation parameter  $\tau^{[1]}$ 
  - Old trick used to solve (continuous) Boltzmann equation<sup>[2]</sup> for low density gases (assuming pairs of particles collide without prior correlation)
  - Most commonly-used form of LBE for simulations

Range of useable lattices expanded<sup>[1]</sup>: square lattices with multiple speeds will now work!

$$\Omega_{ij} \to -\frac{\delta_{ij}}{\tau}$$

No matrix operations now necessary, so even more speed-up

$$f_i(\mathbf{x} + \mathbf{e}_i \Delta t, t + \Delta t) - f_i(\mathbf{x}, t) = -\frac{1}{\tau} \Big( f_i(\mathbf{x}, t) - f_i^{eq} \big( \rho(\mathbf{x}, t), \mathbf{u}(\mathbf{x}, t) \big) \Big)$$



Qian et al., EPL 17 (6), 479–484 (1992) 2.

Bhatnagar et al., Phys Rev 94 (3), 511–525 (1954)

He and Luo, PRE 56 (6), 6811–6817 (1997)

Discretising continuous Boltzmann equation also gets us here<sup>[3]</sup>

### LBE solves most issues with LGCA...

- No ensemble-averaging required afterwards
- No statistical noise: wider viscosity range
- Many more lattice schemes possible, including square-grid multispeed ones, with guaranteed isotropy (Galilean invariance)
- Simpler and cheaper collisions
- Multi-physics and parallelisation still possible
- Still a particle-based method (underneath)
- But ...
  - Floating point numbers instead of integers: no longer unconditionally stable
  - Lattices still usually isothermal/athermal: often cannot provide sufficient particle speeds to obtain definable temperature\* and deal with very compressible flows



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 $f_i(\mathbf{x} + \mathbf{e}_i \Delta t, t + \Delta t) - f_i(\mathbf{x}, t)$  $= -\frac{1}{\tau} \Big( f_i(\mathbf{x}, t) - f_i^{eq} \big( \rho(\mathbf{x}, t), \mathbf{u}(\mathbf{x}, t) \big) \Big)$ 

Local equilibrium based on Maxwell-Boltzmann distribution:

$$f_i^{eq}(\rho, \mathbf{u}) = \frac{\rho}{(2\pi RT)^{D/2}} \exp\left(-\frac{m(\mathbf{e}_i - \mathbf{u})^2}{2k_B T}\right)$$

but usually approximated for low Mach numbers  $(u \ll c_s = \sqrt{3k_B T/m})$  to:

$$f_i^{eq}(\rho, u) = \rho w_i [1 + 3(\mathbf{e}_i \cdot \mathbf{u}) + \frac{9}{2}(\mathbf{e}_i \cdot \mathbf{u})^2 - \frac{3}{2}u^2]$$
  
where  $c_s^2 = \frac{1}{3} \left(\frac{\Delta x}{\Delta t}\right)^2$ ,  $p = \rho c_s^2$  and  
 $\nu = \frac{\mu}{\rho} = \frac{1}{3} \left(\tau - \frac{1}{2}\right) \frac{(\Delta x)^2}{\Delta t}$ 

### Connecting LBE to Navier-Stokes<sup>[1]</sup>

Assume small Knudsen number  $\epsilon$ (molecular mean free path/length scale)

$$f_{i} = f_{i}^{eq} + \epsilon f_{i}^{(1)} + \epsilon^{2} f_{i}^{(2)} + \dots$$
$$\sum_{i} f_{i}^{eq} = \rho$$
$$\sum_{i} \mathbf{e}_{I} f_{i}^{eq} = \rho \mathbf{u}$$
$$\sum_{i} f_{i}^{(k>0)} = \sum_{i} \mathbf{e}_{i} f_{i}^{(k>0)} = 0$$

Length and time scales for convection (1) and diffusion (2) (Chapman-Enskog multiscaling)

$$\frac{\partial}{\partial x} = \epsilon \frac{\partial}{\partial x_1} + \cdots$$
$$\frac{\partial}{\partial t} = \epsilon \frac{\partial}{\partial t_1} + \epsilon^2 \frac{\partial}{\partial t_2} + \cdots$$

Science and Technology Facilities Council Taylor expand LBE about  $\epsilon$ , put in multiscale derivatives, collect terms for  $\epsilon^1$  and  $\epsilon^2$ 

Fast modes 
$$(\epsilon^{1})$$
:  
 $\partial_{t1}f_{i}^{eq} + \mathbf{e}_{i} \cdot \nabla_{1}f_{i}^{eq} = \frac{1}{\tau}f_{i}^{(1)}$   
Slow modes  $(\epsilon^{2})$ :  
 $\partial_{t2}f_{i}^{eq} + \left(1 - \frac{1}{2\tau}\right)(\partial_{t1} + \mathbf{e}_{i} \cdot \nabla_{1})f_{i}^{(1)} = -\frac{1}{\tau}f_{i}^{(2)}$   
Take moments in  
 $\mathbf{e}_{i}$ , sum over  
lattice links *i* and

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{u}) = 0$$
$$\rho \left( \frac{\partial \mathbf{u}}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u} \right) = -\nabla p + \nu \nabla \cdot \left( \nabla (\rho \mathbf{u}) + \left( \nabla (\rho \mathbf{u}) \right)^T \right)$$

Definitions for p,  $c_s$  and vemerge from procedure (results on previous slide)

recombine scales

1. Chen and Doolen, Annu Rev Fluid Mech **30** (1), 329–364 (1998)

# Can start by treating LBE as mesoscopic CFD solver

- Parameterisation mainly top-down: getting lattice spacing  $\Delta x$  and timestep  $\Delta t$
- Use expressions for speed of sound *c<sub>s</sub>* and kinematic viscosity ν with known values for given fluid and a choice of τ (relaxation time)
  - $\tau$  cannot be too close to  $\frac{1}{2}$  (zero viscosity) or calculation will become numerically unstable



Example – water at room temperature (298 K):

Setting  $\tau = 1$   $c_s = 1498 \text{ m s}^{-1}$   $\nu = 10^{-6} \text{ m}^2 \text{ s}^{-1}$   $\Delta x \approx 2.31 \times 10^{-8} \text{ m} (23.1 \text{ nm}, 0.0231 \text{ µm})$  $\Delta t \approx 8.91 \times 10^{-13} \text{ s} (0.891 \text{ ps})$ 

- Density ρ is a relatively free parameter, but usually set to 1 to maximise available calculation precision
- No need to define exactly what particles represent on lattice: we can use LBE for liquids as well as gases

### Using LBE for mesoscale modelling Boundary conditions

- Finding 'missing' values of  $f_i$  re-entering grid
- Rules can be simple/intuitive and applied at any grid point
  - Bounce-back reflection of  $f_i$  to give no-slip boundary<sup>[1]</sup>
  - 'Equilibrium forcing': replacing  $f_i$  with  $f_i^{eq}(\rho_w, \mathbf{u}_w)$
- More sophisticated boundary conditions available to set  $\rho_w$  and/or  $\mathbf{u}_w^{[2,3,4]}$
- Can also devise moving boundaries for suspended
   objects<sup>[5,6]</sup>
  - 1. Lavallée *et al.*, *Physica D* **47** (1–2), 233–240 (1991)
  - 2. Inamuro *et al.*, *Phys Fluids* **7** (12), 2928–2930 (1995)
  - 3. Zou and He, *Phys Fluids* **9** (6), 1591–1598 (1997)
  - 4. Halliday et al., J Phys A **35** (12), L157–L166 (2002)
  - 5. Ladd, *J Fluid Mech* **271**, 285–309 (1994)
  - 6. Feng and Michaelides, J Comput Phys 195 (2), 602-628 (2004)





 $f_{-i}(\mathbf{x}_w, t + \Delta t) = f_i(\mathbf{x}_w, t) + 2w_i \rho \mathbf{U}_p \cdot \mathbf{e}_i$ 

Can work out force acting on object from  $f_i$  at grid points surrounding boundary and determine its motion (translational and rotational)



### Interactions between multiple fluids and phases

- Can add a force **F** acting on a fluid into collisions by:
  - Modifying fluid velocity used in  $f_i^{eq}$ <sup>[1]</sup>
  - Adding a forcing term to the collision<sup>[2]</sup>
  - A combination of both<sup>[3]</sup>
- Possible to devise forces based on gradients in densities (or similar properties)
  - Easy to calculate gradients of a property  $\phi$  accurately using a lattice stencil, e.g.

$$c_s^2 \nabla \phi(\mathbf{x}) \approx \sum_i w_i \phi(\mathbf{x} + \mathbf{e}_i \Delta t) \mathbf{e}_i$$

• Multiple fluids modelled on individual lattice grids (occupying same spatial locations) with their own distribution functions



- 1. Martys and Chen, PRE 53 (1), 743–750 (1996)
- 2. Kupershtokh and Medvedev, J Electrostat 64 (7–9), 581–585 (2006)
- 3. Guo et al., PRE 65 (4), 046308 (2002)

### Chromodynamics<sup>[1]</sup>

- Top-down approach: designed to get correct hydrodynamics and interfacial tension between fluids, originally used for multiplecomponent lattice gases<sup>[2]</sup>
- Calculate densities at each lattice point ( $\rho_a = \sum_i f_i^a$ ) and then a 'phase index':

$$\rho^N = \frac{\rho_a - \rho_b}{\rho_a + \rho_b}$$

- Forces based on phase index gradients and interfacial tensions
- Fluids collided together 'achromatically' before being re-segregated
  - Combining distribution functions for collision ensures correct hydrodynamics
  - Segregation produces diffuse interfaces, but avoids 'pinning' of drops to lattice and allows larger interfacial tensions



Pure component *a*: Pure component *b*: Interface:





1. Lishchuk et al., PRE 67 (3), 036701 (2003)

2. Gunstensen *et al.*, *PRA* **43** (8), 4320–4327 (1991)

### **Pseudopotentials**<sup>[1]</sup>

- Makes better connection to microscopic interactions, can be used for single fluid (multiple phases) as well as multiple fluids
- Define 'pseudopotential'  $\psi(\rho)$  as function of fluid density for each grid point
- Force on fluid(s) based on gradient(s) of pseudopotential(s):

$$\mathbf{F}^{a} = -\psi^{a}(\mathbf{x}) \sum_{b} g_{ab} \sum_{i} w_{i} \psi^{b}(\mathbf{x} + \mathbf{e}_{i} \Delta t) \mathbf{e}_{i}$$

• Form of pseudopotential determines equation of state for fluid(s) ...

$$p = \rho c_s^2 + \frac{1}{2}g c_s^2 \psi^2$$

• ... or vice versa!<sup>[2]</sup>





Most commonly-used LBE algorithm for multiple phases or fluids, but not without its flaws:

- Spurious flows generated at fluid/phase interfaces
- Often not thermodynamically consistent
- No interfacial tension control
- 1. Shan and Chen, *PRE* **47** (3), 1815–1819 (1993)
- 2. Peng and Schaefer, *Phys Fluid* **18**, 042102 (2006)

### Using LBE for mesoscale modelling Free-energy LBE<sup>[1]</sup>

• Modifies  $f_i^{eq}$  by adding in bulk pressure and interfacial terms from density gradients:

$$f_i^{eq}(\rho, u) = \rho w_i^{00} + w_i \left[ \rho(\mathbf{e}_i \cdot \boldsymbol{u}) + \frac{3}{2}\rho(\mathbf{e}_i \cdot \boldsymbol{u})^2 - \frac{1}{2}\rho u^2 + f(\lambda, \mathbf{e}_i, \nabla \rho) \right] + w_i^p P_0 - w_i^t \kappa \nabla^2 \rho + \kappa \sum_{\alpha, \beta} w_i^{\alpha\beta} \partial_\alpha \rho \partial_\beta \rho$$

- Imposes equation of state via bulk pressure  $P_0(\rho, T)$ , correct and consistent thermodynamics
- No additional force added to fluid, but need EOS-dependent correction to ensure Galilean invariance
- Density gradients calculated using accurate grid stencils<sup>[2]</sup>
- Can apply similar (but simpler) local equilibrium distribution function for concentration of two-fluid mixtures, setting chemical potential instead of bulk pressure
  - Harder to apply for more than two fluids





1. Swift et al., PRL 75 (5), L157–L166 (2002)

2. Pooley and Furtado, *PRE* 77, 046702 (2008)

### Solute diffusion, reactions and heat transfer

- We can change  $f_i^{eq}$  to represent exactly incompressible liquids<sup>[1]</sup> or truncate to linear velocity terms for low-speed diffusion
- Diffusive solutes can be modelled as additional fluids, but coupled to solvent fluid motion through their own  $f_i^{eq}$ <sup>[2]</sup>
  - Source terms can be added into collisions to represent reaction kinetics<sup>[3]</sup> and catalytic reactions can be applied at boundaries<sup>[4]</sup>
- Same approach can be used for heat transfer
  - Simulations of heat convection possible with fluid interactions<sup>[2]</sup> and combustion with reactions/solid comsumption<sup>[5]</sup>





- 1. He and Luo, *J Stat Phys* **88** (3–4), 927–944 (1997)
- 2. Inamuro et al., J Comput Phys 179 (1), 201–215 (2002)
- 3. Ponce Dawson et al., J Chem Phys 98 (2), 046308 (1993)
- 4. Arcidiacono et al., PRE 78 (4), 046711 (2008)
- 5. Liu et al., Combust Flame **211**, 325–336 (2020)

# Rheology, turbulence and viscoelasticity

- To apply a rheological model to a fluid, calculate ν
   (τ) at each grid point based on its local shear rate
  - Shear rate can be obtained either from velocity gradients or locally from non-equilibrium part of momentum flux tensor (second moment of  $f_i$ )<sup>[1]</sup>
- Turbulence can be applied this way, either with two-equation<sup>[2]</sup> or Large Eddy Simulation models<sup>[3]</sup>
- Viscoelasticity also available using similar approach<sup>[4]</sup>



- 1. Boyd *et al.*, *J Phys A* **39** (46), 14241–14247 (2006)
- 2. Filippova et al., J Comput Phys 170 (2), 812-829 (2001)
- 3. Yu et al., Comput Fluids 35 (8), 957–965 (2006)
- 4. Malaspinas et al., J Nonnewton Fluid Mech 165 (23–24), 1637–1653 (2010)





### LBE challenges: Collisions

- Simple collisions (BGK) with single relaxation time  $\tau$ 
  - Numerical instabilities for very low viscosities and higher flow speeds
  - No direct control over bulk viscosity
- May need to resort to more complex collisions with additional parameters
  - Two-Relaxation-Time (TRT) collisions<sup>[1]</sup> work on conjugate  $f_i$  values: use additional relaxation time to help reduce numerical instabilities
  - Multiple Relaxation Time (MRT) collisions<sup>[2]</sup> work on additional f<sub>i</sub> moments: bulk viscosity controlled by an additional relaxation time, other relaxation times help eliminate higher-order non-hydrodynamic moments
  - Cascaded LBE (CLBE) collisions<sup>[3,4]</sup> work on central moments (moments about  $\mathbf{e}_i \mathbf{u}$ ), also uses higher-order  $f_i^{eq}$  greater stability than other MRT collisions



- 1. Ginzburg et al., Commun Comput Phys 3 (2), 427–478 (2008)
- 2. Lallemand and Luo, *PRE* **61** (6), 6546–6562 (2000)
  - 3. Geier et al., PRE 73 (6), 066705 (2006)
- 4. Lycett-Brown and Luo, Comput Math Appl 67 (2), 350-362 (2014)

### LBE challenges: Phase/fluid interfaces

- Spurious microcurrents generated at interfaces between fluids or phases
  - Particular problem for large contrasts in density between phases/fluids and/or higher flow speeds
  - Multi-component LBE schemes available to work with larger density contrasts (e.g. gas/liquid)<sup>[1]</sup>
  - Stencils to calculate gradients more accurately<sup>[2]</sup>
  - Use of multiple-relaxation-time collisions can help
- Diffuse interfaces: a few lattice points wide<sup>[3]</sup>
- Interfacial tension not directly controlled in pseudopotential model<sup>[4]</sup>



- Lishchuk *et al.*, *PRE* **77** (3), 036702 (2008)
   Pooley and Furtado, *PRE* **77** (4), 046702 (2008)
- 3. Spencer and Halliday, *PRE* 88 (6), 063305 (2013)
- 4. Li and Luo, *PRE* 88 (5), 053307 (2013)











### LBE challenges: Thermohydrodynamics

- Most LBE lattices use too few links to define temperature
  - Can couple heat transfer effects to fluid flow using separate lattice (but does not solve compressibility issue)<sup>225</sup>
- Low density (rarefied) gas flows with larger ε: going beyond Navier-Stokes
  - Slip flows can occur at boundaries even for low Ma<sup>[1]</sup>
  - Lattice schemes with many more speeds<sup>[2]</sup>
- Entropic LBE schemes: devise  $f_i^{eq}$  to additionally satisfy Boltzmann *H* theorem<sup>[3]</sup>
  - Can enable more compressible flows than normally available (including supersonic)



- 1. Zhang et al., PRE 71 (4), 047702 (2005)
- 2. Shan et al., *J Fluid Mech* **550**, 413–441 (2006)
- 3. Karlin *et al.*, *EPL* **47** (2), 182–188 (1999)
- 4. Frapolli et al., Entropy 22 (3), 370 (2020)





Entropic LBE simulations of supersonic flow around airfoil (top) and vortex interacting with shock wave (bottom)<sup>[4]</sup>

### More information about LBE

• Knowledge Centre in DL\_Software Digital Guide:

#### https://dl-sdg.github.io/RESOURCES/knowledge.html

 Includes pages on LBE and further details (how to set scales in LBE calculations, more details on multiple-phase/fluid algorithms, boundary conditions etc.)





### LBE in DL\_Software: DL\_MESO\_LBE

'Let us think the unthinkable, let us do the undoable, let us prepare to grapple with the ineffable itself, and see if we may not eff it after all.'

- Douglas Adams, *Dirk Gently's Holistic Detective Agency* 

### **DL\_MESO**

# General purpose mesoscopic simulation software package

- MPI domain-decomposed codes with optional OpenMP multithreading for:
  - Dissipative Particle Dynamics (DPD)
  - Lattice Boltzmann Equation (LBE)
- Created in 2004 as CCP5 flagship project: www.ccp5.ac.uk/DL\_MESO
  - Development currently funded under CoSeC for UKCOMES
- Current version (2.7): released December 2018, 700+ registered academic users





### **DL\_MESO**

# General purpose mesoscopic simulation software package

- Two articles available in Molecular Simulation on DL\_MESO
- Cited ~200 times over past 12 years
- Available under dual licence:
  - Free for academic users
  - Annual subscription for commercial users



Molecular Simulation, 2013 Vol. 39, No. 10, 796-821, http://dx.doi.org/10.1080/08927022.2013.772297



#### DL\_MESO: highly scalable mesoscale simulations

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DL\_MESO is a parallel mesoscale simulation package capable of dissipative particle dynamics and the lattice Boltzmann equation method. It has been developed at Daresbury Laboratory for the United Kingdom Collaborative Computational Project known as CCP5. Capable of addressing industrially relevant tasks, but written to support academic research, it has a wide range of applications and scales to thousands of processors on high-performance computing platforms yet runs efficiently on smaller commodity clusters and single processor personal computers. This article serves as a guide to a variety of users, describing the functionality, performance and structure of this simulation package. Representative examples highlighting the capabilities of DL\_MESO are given for each of the two methodologies available. Future directions for the package are discussed towards the end of the article.

Keywords: DL\_MESO; dissipative particle dynamics; lattice Boltzmann; mesoscale simulation

#### 1. Introduction

Computer simulations of condensed matter are frequently performed using either atomistic methods, e.g. molecular dynamics (MD), or continuum methods based on the Navier–Stokes equation. Although these are very successful in their own domains, many systems exist in which large-scale structures appear but require some vestige of atomistic detail to capture the essential physics. Mesoscale m

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scales, which are important in the context of consumer interaction with formulated products.

The DL\_MESO project originated as part of the United Kingdom Collaborative Computational Project for the Computer Simulation of Condensed Phases, known as CCP5 (www.ccp5.ac.uk). The objective was to develop a comprehensive mesoscale modelling package capable of bridging the gap between atomistic and continuum



Check for updates

#### DL\_MESO\_DPD: development and use of mesoscale modelling software

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#### ABSTRACT

DL\_MESO is a highly-scalable general purpose software package for mesoscale modelling. Created and developed at Daresbury Laboratory for the UK Collaborative Computational Project CCP5, it was intended to be a companion package to the flagship molecular dynamics code DL\_POLY. One of DL\_MESO component codes, DL\_MESO\_DPD, is based on dissipative particle dynamics, a mesoscale modelling technique with many similarities to classical molecular dynamics. While this code and DL\_POLY were created with different applications in mind, they share a significant amount of functionality and development history. This article gives an overview on how DL\_MESO\_DPD has been developed, including its shared history with DL\_POLY and information on its current performance, and a selection of applications for which the code has been used.

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#### KEYWORDS

DL\_MESO; mesoscale modelling; dissipative particle dynamics; software development; DL\_POLY

#### 1. Introduction

Mesoscale modelling methods fit into a gap between atomistic and continuum methods of modelling materials, addressing intermediate length and timescales (10–1000 nm, 1 ns–10 ms) by including some vestige of atomistic detail to capture both essential microscopic physics and large-scale structural effects. As well as being of interest to the academic community, mesoscopic simulation methods are of great interest to industry and many of their key developments

The DPD part of the DL\_MESO package, DL\_MESO\_DPD, has a lot in common with DL\_POLY. Both programs and their included methodologies (DPD and MD) depend upon calculating interaction forces acting on particles and integrating those forces over time to determine their motion. From the user's point of view, both programs use similar input file formats which are – to at least some degree – interchangeable and compatible. DL\_MESO\_DPD and DL\_POLY\_4 also share the same parallelisation strategy of domain decomposition with link-cell

### DL\_MESO\_LBE

### **Flexible LBE solver**

- Solves discretised Boltzmann equation on regular lattice(s)
- Parallelised mainly by domain decomposition using MPI, but can also use OpenMP multithreading – both share grid points among cores/threads
- Can specify boundary conditions for walls/obstacles, initial conditions and fluid interaction and properties using input files
- Utilities provided to create initial system states, collect up output files created in parallel for visualisation and analysis
  - Also possible to write scripts for simulation setup and post-processing



### **DL\_MESO\_LBE: current functionalities**

#### Square lattice grids (two- and three-dimensional)

• D2Q9, D3Q15, D3Q19, D3Q27

#### Multiple fluids available

- Weakly compressible or fully incompressible fluids
- Can couple in solute mass transfers and/or heat transfers

#### **Collision schemes**

- BGK single relaxation time (default)<sup>[1]</sup>
- Two Relaxation Time (TRT)<sup>[2]</sup>
- Multiple Relaxation Time (MRT)<sup>[3]</sup>
- Cascaded Lattice Boltzmann (CLBE)<sup>[4]</sup>



Different schemes to apply external and interfacial forces available for all collision schemes

- 1. Qian et al., EPL 17 (6), 479–484 (1992)
- 2. Ginzburg et al., Comm Comp Phys 3 (2), 427–478 (2008)
- 3. Lallemand and Luo, PRE 61 (6), 6546–6562 (2000)
- 4. Geier et al., PRE 73, 066705 (2006)

### **DL\_MESO\_LBE: current functionalities**

#### **Boundary conditions**

- No-slip bounce-back conditions (simplest)
- Outflows
- Fixed velocity or density (Zou/He<sup>[1]</sup>, Inamuro<sup>[2]</sup>, regularised<sup>[3]</sup>, kinetic<sup>[4]</sup> BCs)
- Fixed concentration and temperature (Zou/He, Inamuro)
- Surface wetting available

#### **Mesoscopic interactions**

- Shan/Chen pseudopotential for any number of fluids
  - Cubic and hard-sphere equations of state
- Lishchuk (chromodynamic) continuum-based for two or more fluids
- Swift free-energy for one/two fluid systems
  - Cubic and hard-sphere equations of state
- 1. Zou and He, Phys Fluid 9 (6), 1591–1598 (1997)
- 2. Inamuro et al., Phys Fluid 7 (12), 2928–2930 (1995)
- 3. Latt et al., PRE 77, 056703 (2008)
- 4. Ansumali and Karlin, *PRE* **66**, 026311 (2002)

### **DL\_MESO\_LBE: current functionalities**

#### **Non-Newtonian rheological models**

- Uses momentum flux tensors to calculate shear rates locally, used as inputs for rheological models
- Power law fluids, Bingham/Herschel-Bulkley plastics, Casson and Carreau-Yasuda fluids

#### **Parallel file-writing**

- Options to gather data along Cartesian axes to reduce number of output files produced (normally one file per core per simulation snapshot)
- Gathering in all three directions invokes MPI-IO generates one file per simulation snapshot



- 1. Zou and He, *Phys Fluid* **9** (6), 1591–1598 (1997)
- 2. Inamuro et al., Phys Fluid 7 (12), 2928–2930 (1995)
- 3. Latt *et al.*, *PRE* **77**, 056703 (2008)
- 4. Ansumali and Karlin, PRE 66, 026311 (2002)

### Future DL\_MESO\_LBE functionalities

- Non-spherical drops (fluid-filled vesicles)<sup>[1]</sup>
  - Adaptation of chromodynamic (Lishchuk) algorithm to include conservation of membrane area and bending
  - Useful for modelling e.g. red blood cells in haemodynamic flows
- Suspended solid object dynamics<sup>[2]</sup>
  - Calculating forces on objects and integrating to determine motion
  - Opens up coupling between MD/DPD and LBE
- Adaptive mesh LBE
  - Creating finer grids (smaller Δx) close to interfaces for greater accuracy, e.g. curved boundaries
  - Coupling together fine and coarse grids



- 1. Spendlove et al., Phys Fluid **33**, 077110 (2021)
- 2. Ladd, J Fluid Mech 271, 285–309 (1994)

### **LBE** units

#### Based upon lattice (both set to 1 inside DL\_MESO\_LBE)

- Grid spacing  $\Delta x$
- Time step  $\Delta t$

#### Relaxation time $\tau$ and 'real-life' quantities can be used to find lattice units

- Speed of sound  $c_s = \frac{1}{\sqrt{3}} \frac{\Delta x}{\Delta t}$
- Kinematic viscosity  $\nu = \frac{\mu}{\rho} = \frac{1}{3} \left( \tau \frac{1}{2} \right) \frac{(\Delta x)^2}{\Delta t}$

#### **Density in LBE calculation a 'free' parameter**

- Multiplier in local equilibrium distribution function  $f_i^{eq}$
- Usually set to (or around) 1 to maximise calculational accuracy

### **DL\_MESO\_LBE input/output files**

#### **Input files**

- Ibin.sys
   Simulation parameters/settings
- Ibin.spa Boundary conditions
  - Ibin.init Initial conditions/configuration

#### **Output files**

- Screen System diagnostics (can re-direct to file)
- Ibout\*.vts
   Plottable results in VTK format (default: other options available)
- Ibout.dump Simulation restart capability
- Ibout.info Simulation information for post-processing after parallel run
   Ibout.ext

### Ibin.sys – Simulation parameters/settings

- Mandatory simulation control and system definition
  - Similar to CONTROL and FIELD in DL\_MESO\_DPD/DL\_POLY
- Each line is a key phrase followed by a number/word
  - First 10 in bold **must** be specified
  - Parameters not required can be omitted
  - No specific line order
- Can specify
  - Initial/boundary velocities, fluid densities, solute concentrations, temperature
  - Mesoscopic interaction and rheological models/parameters



Chapter 6 in DL\_MESO User Manual gives full list of available key phrases

space_dimension	2				
discrete_speed	9				
number_of_fluid	2				
number_of_solute	0				
temperature_scalar	0				
phase_field	0				
grid_number_x	150				
grid_number_y	50				
grid_number_z	1				
domain_boundary_width	1				
incompressible_fluids	0				
collision_type	BGKGuo				
interaction_type	Lishchuk				
output_format	VTK				
total_step	40000				
equilibration_step	5000				
save_span	500				
noise_intensity	0.0				
sound_speed	540.0				
kinetic_viscosity	0.001				
relaxation_fluid_0	1.0				
relaxation_fluid_1	1.0				
bulk_relaxation_fluid_0	1.0				
bulk_relaxation_fluid_1	1.0				
interaction_0	0.0				
interaction_1	0.0034				
interaction_2	0.0034				
interaction_3	0.0				
segregation	1.4				

### Ibin.spa – Boundary conditions

	_	_	_	_	_	_	_	_	_	_	_	_
1	3	1		4	9		0		1	4	9	
1	3	2		4	9		0		1	4	9	
1	3	3		4	9		0		1	4	9	
1	3	4		4	9		0		1	4	9	
1	3	5		4	9		0		1	4	9	
1	3	6		4	9		0		1	4	9	
1	3	7		4	9		0		1	4	9	
1	3	8		4	9		0		1	4	9	
1	3	9		4	9		0		1	4	9	
1	4	0		4	9		0		1	4	9	
1	4	1		4	9		0		1	4	9	
1	4	2		4	9		0		1	4	9	
1	4	3		4	9		0		1	4	9	
1	4	4		4	9		0		1	4	9	
1	4	5		4	9		0		1	4	9	
1	4	6		4	9		0		1	4	9	
1	4	7		4	9		0		1	4	9	
0		4	9		0		1	4	9			
1	4	9		4	9		0		1	4	9	
1		0		0		1	3					
1	4	8		0		0		1	3			
2		0		0		1	3					
3		0		0		1	3					
4		0		0		1	3					
5		0		0		1	3					
6		0		0		1	3					

- Mandatory definition of boundary conditions
  - A blank file still needed even if all BCs are periodic
- Four numbers per line representing a grid point

```
x y z [BC code]
```

- BC codes listed in chapter 6 of DL\_MESO user manual
  - Blank sites, on-grid and mid-grid bounce-back
  - Outflow boundaries
  - Fixed velocity/fluid density, solute concentration, temperature
- BC codes for outflow and fixed velocity/density/concentration/ temperature boundaries are directional, blank sites and bounceback are not
- Only need to specify BC code for a grid point if its boundary condition is <u>not</u> periodic
- Easiest to create **Ibin.spa** using DL\_MESO GUI or a script

### Ibin.init – Initial conditions

26 15 0 0.0 0.0 0.0 0.0 2.0 26 16 0 0.0 0.0 0.0 0.0 2.0 26 17 0 0.0 0.0 0.0 0.0 2.0 26 18 0 0.0 0.0 0.0 0.0 2.0 26 19 0 0.0 0.0 0.0 0.0 2.0 26 20 0 0.0 0.0 0.0 0.0 2.0 26 21 0 0.0 0.0 0.0 0.0 2.0 26 22 0 0.0 0.0 0.0 0.0 2.0 26 23 0 0.0 0.0 0.0 0.0 2.0 26 24 0 0.0 0.0 0.0 0.0 2.0 26 25 0 0.0 0.0 0.0 0.0 2.0 27 13 0 0.0 0.0 0.0 0.0 2.0 27 14 0 0.0 0.0 0.0 0.0 2.0 27 15 0 0.0 0.0 0.0 0.0 2.0 27 16 0 0.0 0.0 0.0 0.0 2.0 17 0 0.0 0.0 0.0 0.0 2.0 18 0 0.0 0.0 0.0 0.0 2.0 27 19 0 0.0 0.0 0.0 0.0 2.0 27 20 0 0.0 0.0 0.0 0.0 2.0 27 21 0 0.0 0.0 0.0 0.0 2.0 27 22 0 0.0 0.0 0.0 0.0 2.0 27 23 0 0.0 0.0 0.0 0.0 2.0 27 24 0 0.0 0.0 0.0 0.0 2.0 27 25 0 0.0 0.0 0.0 0.0 2.0 27 26 0 0.0 0.0 0.0 0.0 2.0 27 27 0 0.0 0.0 0.0 0.0 2.0 28 12 0 0.0 0.0 0.0 0.0 2.0 28 13 0 0.0 0.0 0.0 0.0 2.0 28 14 0 0.0 0.0 0.0 0.0 2.0

- Optional simulation initialisation
- Each line specifies grid point coordinate and properties for that grid point: velocity, fluid densities, solute concentrations and temperature
  - Similar format to **lbin.spa**:

 $\mathbf{x} \mathbf{y} \mathbf{z} \mathbf{v}_{\mathbf{x}} \mathbf{v}_{\mathbf{y}} \mathbf{v}_{\mathbf{z}} \mathbf{\rho}_{0} \mathbf{\rho}_{1} \dots$ 

- All properties for system <u>must</u> be specified in each line
- Values for specified points used in place of default system-wide values given in lbin.sys
- Utilities supplied with DL\_MESO can be used to create lbin.init
  - Ibeinitcreate.cpp (init.exe) adds spherical/circular drops for new simulations
  - **Ibedumptoinit.cpp** (dump\_to\_init.exe) creates lbin.init from lbout.dump restart file
- Also possible to devise scripts to write this file

### **Screen/standard output – Information**

- DL\_MESO\_LBE outputs information about calculation to screen or standard output
  - Can pipe this to a file
  - Similar to DL\_MESO\_DPD/DL\_POLY OUTPUT
- Starts with some information about simulation (e.g. grid size, numbers of fluids)
- Writes out masses of fluids and system momentum at user-specified intervals
  - Useful as diagnostic for LBE calculation, e.g. checking progress towards steady-state
- At end of calculation, prints timing information and efficiency measure: Millions of Lattice Updates Per Second (MLUPS)

Welcome to DL\_MESO\_LBE

1.31106

```
UKRI STFC/CCP5 Program Library Package
UKRI STFC Daresbury Laboratory Lattice Boltzmann Equation program
DL MESO version 2.7 rev 01, January 2019
Authors:
              R. S. Qin & M. A. Seaton
Contributors: J. Meng, L. Grinberg & M. A. Johnston
Copyright (c) 2019 UKRI STFC Daresbury Laboratory, United Kingdom
Any comments or queries, please contact Dr. M. A. Seaton at:
   Email: michael.seaton@stfc.ac.uk
   Tel: +44 (0)1925 603317
Program started at: Sun Jan 6 07:58:28 2019
armed with D2O9 LB model
2-dimensional system with grid size: nx=150, ny=50
boundary width = 0
system includes 2 compressible phases, 0 solute and no temperature scalar
D2Q9 lattice Boltzmann model is defined using BGK collisions
with Guo forcing and Lishchuk interactions
Running with 8 threads
fluid 0 :
                         simple Newtonian fluid
                         relaxation time
                                                    = 1
fluid 1 :
                         simple Newtonian fluid
                         relaxation time
                                                    = 1
interaction parameters: fluid i fluid j
                                                                beta ij
                                                    g ij
                                                  0.0034
                                                                    1 4
                               0
wall interaction parameters:
                                  fluid i
                                                    g wi
             MASS: total = 14700, fluid 0 = 13306, fluid 1 = 1394
0
0.0101249
             MOMENTUM: x: -2.04003e-13, y: -2.04003e-13
500
             MASS: total = 14700, fluid 0 = 13306, fluid 1 = 1394
0.665448
            MOMENTUM: x: 5.97018e-12, y: 0.0313264
1000
             MASS: total = 14700, fluid 0 = 13306, fluid 1 = 1394
```

MOMENTUM: x: 5.75585e-12, y: 0.0357875

### lbout\*.vts - Simulation snapshots/trajectory

- Snapshots of simulation at user-specified intervals as trajectory
  - Similar to HISTORY files in DL\_MESO\_DPD and DL\_POLY
  - Includes macroscopic properties for each grid point (velocities, densities, concentrations, temperatures) and boundary condition
- Default format: XML-based Structured Grid VTK
  - Legacy VTK and Plot3D formats also available
  - Binary and text variants available
  - All formats readable in ParaView
- At least one file produced per snapshot: by default, each core writes its own files (to be gathered afterwards)
  - Strongly recommend using data gathering options in Ibin.sys when running in parallel to reduce number of files



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### **Ibout.dump – Simulation restart**

- Binary file with distribution functions (and local relaxation times when using rheological models) for all lattice points
  - Similar to export/REVCON and REVIVE for DL\_MESO\_DPD/DL\_POLY
  - DL\_MESO\_LBE can read this file to resume a previous calculation
  - Ibedumpvtk.cpp (dump\_to\_vtk.exe) utility creates VTK file of simulation state in Ibout.dump for visualisation (and simulation diagnostics)
  - Ibedumptoinit.cpp (dump\_to\_init.exe) utility creates Ibin.init as initial state for new calculation

More information available in DL\_Software Digital Guide about DL\_MESO\_LBE files and formats:

https://dl-sdg.github.io/RESOURCES/TUTORIALS/dlm\_3.html

and utilities:



https://dl-sdg.github.io/RESOURCES/TUTORIALS/dlm\_6.html

### **LBE Tutorial Exercises**

- Three available in DL\_Software Digital Guide, all designed to work with DL\_MESO\_LBE
- Phase separation and equations of state
  - Setting up Shan/Chen pseudopotential interactions to model vapour-liquid coexistence from a cubic equation of state
- Drop flows
  - Using chromodynamics to model drops undergoing linear shear, examining effects of shear rate and interfacial tension



### **LBE Tutorial Exercises**

- Pressure-driven flows and Kármán vortex streets
  - Applying channel flows with obstacles
- All three exercises use DL\_MESO\_LBE with input files supplied, some Python scripts provided for data analysis (Phase separation) and boundary condition building (Pressure-driven flows), and ParaView for visualisation

#### https://dl-sdg.github.io/RESOURCES/EXERCISES/Exercises.html





# Questions?



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# Thank you

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