

DARESBUURY LABORATORY

INFORMATION QUARTERLY

for

COMPUTER SIMULATION OF CONDENSED PHASES

An informal Newsletter associated with Collaborative Computational Project No. 5
on Molecular Dynamics, Monte Carlo and Lattice Simulations of Condensed Phases.

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Editorial.

It is probably part of the natural cycle of things that it occasionally becomes necessary to appeal for contributions for the CCP5 Newsletter. Such is the case now. This does not however indicate that there is a diminishing interest in the newsletter itself, as rarely a day goes by when we are not asked to add more names to our mailing list. It is therefore important to remind our readers that the newsletter is very dependent on their enthusiasm if it is to keep going.

The CCP5 newsletter is an interesting publication in a number of ways: it published regularly, articles appear quickly (within three months or less), it has an international readership and it is informal enough to allow the publication of novel or speculative ideas. Added to these is the fact that there are surprisingly few publications that discuss the computational details of simulation methods warts and all. For these reasons it should be well worth supporting.

So to encourage our readership we announce that we would be pleased to receive contributions from anyone active in the field of computer simulation of condensed phases, anywhere in the world! We would also be particularly pleased to publish material from postgraduate researchers, who often have some original ideas to contribute to the computational art. We await your contributions!

In view of this appeal it is a special pleasure to thank the contributors to this month's issue. Long may they prosper!

Contributors:

D. Brown	Department of Chemistry, U.M.I.S.T., Sackville St., Manchester M60 1QD.
C.R.A. Catlow R.A Jackson	Department of Chemistry, University of Keele, Keele, Staffs. ST5 5BG.
J.E. Quinn	DAP Support Unit, Queen Mary College, Mile End Road, London E1 4NS.
W. Smith	Theory and Computational Science Division, S.E.R.C. Daresbury

Laboratory, Daresbury, Warrington
WA4 4AD.

L. Woodcock

Department of Chemical Engineering,
The University of Bradford,
Bradford, Yorkshire BD7 1DP.

General News.

a) CCP5 is to organise a meeting entitled "Industrial Applications of Computer Simulation". The meeting will take place at the Royal Institution in London from 11th. to 13th. January 1988. The purpose of the meeting is to discuss the application of the simulation methods associated with CCP5 (i.e. molecular dynamics, Monte Carlo, energy minimisation etc.) to industrially significant problems. The organisers of the meeting will be Prof. C.R.A. Catlow and Dr. N. Quirke. Readers interested in taking part in this meeting should write to: Dr. N. Quirke, B.P. Research Centre, Chertsey Road, Sunbury-on-Thames, Middlesex TW16 7LN.

Further news of this meeting will be announced later.

b) CCP5 is also to organise a meeting on the subject of bio-organic applications of computer simulation, which provisionally, is scheduled for Easter 1988. More details of this will be revealed as the organisation develops.

c) At the last meeting of the CCP5 Steering Committee which took place on 9th. January the composition of the Executive Committee was changed in accordance with the usual procedure. Professor C.R.A. Catlow was elected as the new CCP5 Chairman to replace Dr. J.H.R. Clarke. In addition Drs. R. Lynden-Bell, S. Parker and M.P. Allen were elected to the Executive Committee. The outgoing members being Drs. N. Quirke and P.A. Madden. In addition to these changes; Dr. M. Leslie has taken the responsibility of Secretary to CCP5 and Dr. W. Smith will now handle the CCP5 Program Library. The composition of the Executive Committee is thus:

Professor C.R.A. Catlow (Keele University and Daresbury)-
Chairman,

Dr. M.P. Allen (Bristol University)

Dr. D.M. Heyes (Royal Holloway and Bedford New College)

Dr. M. Leslie (Daresbury) - Secretary

Dr. R. Lynden-Bell (Cambridge University)

Dr. S. Parker (Bath University)

Dr. W. Smith (Daresbury) - Program Librarian

This would seem to be the place to thank the past members of the Executive Committee for their work on behalf of the CCP5 community. In particular it is appropriate to thank Julian Clarke, under whose chairmanship for the past three years CCP5 has been extremely successful.

d) The CCP5 Executive Committee wish to announce to all U.K. participants in CCP5 that it has set aside a part of the CCP5 budget for the purpose of assisting collaborative work in the U.K. simulation community. These funds are to assist with expenses for travel between collaborating centres. So far, CCP5 has agreed to assist the following collaborating groups:

J. Clarke and D. Tildesley (Transport Properties),
 S. Parker and G. Price (Lattice Simulations),
 M. Allen and G. Evans (Hard Ellipsoids),
 K. Singer and W. Smith (Quantum Simulations).
 Other groups who wish to be considered for this assistance should write to the CCP5 Chairman, Prof. C.R.A. Catlow, Department of Chemistry, University of Keele, Keele, Staffs. ST5 5BG.

e) The CCP5 Executive Committee also wishes to remind U.K. participants that it has an annual allocation of Cray time at each of the centres: London (Cray 1s) and Rutherford (Cray XMP-48), which is available for the development of simulation programs prior to a grant allocation. At present CCP5 is allocated 15 hours a year at London and 5 hours quarterly at Rutherford. Readers who wish to use some of this allocation should write to the CCP5 Secretary, Dr. M. Leslie, TCS Division, SERC Daresbury Laboratory, Daresbury, Warrington WA4 4AD.

f) The Rutherford and Appleton Computer Centre Cray XMP/48 is now up and running. It is currently accessible to a few privileged users, who are running jobs to test the machine in a working environment before it becomes available to the wider academic community in April.

In time it is expected that the CCP5 Program Library will be modified to run on the XMP and users who would like to see any specific programs converted immediately should contact the CCP5 Program Librarian, Dr. W. Smith.

g) The University of Manchester Regional Computing Centre has described the new VM/CMS service which they will be running soon on the Amdahl 5890/300. It was originally intended to run the three domains of the Multiple Domain Facility to deal with: (i) the service currently supplied by the Amdahl 470 V/8; (ii) a national VM/CMS service and (iii) a Manchester local service. However only two domains are possible with the current hardware and so the last two services will run initially on one domain. Some disruption

of the existing service is expected 20-23 March to reconnect the peripherals of the 470 V/8 to the 5890/300, followed by acceptance tests. The experimental VM/CMS service is expected to be available in early May. A useful description of the VM/CMS service appears in the March UMRCC Newsletter. An important aspect of the operation of all the national computing centres (London, Manchester and Rutherford-Appleton) will be the new Peer Review system for the allocation of computing resources. Everyone involved in scientific computing in the UK should become familiar with this. The UMRCC newsletter for March 1987 and the ULCC newsletter for February 1987 both carry details of the agreed procedure and are thus recommended reading.

h) The University of London Computing Centre has now upgraded its front end computer to an Amdahl 5890. The second phase of this upgrade will involve the introduction of new disc storage devices (Amdahl 6380 disc drives, which are 3380 compatible). During this phase both 3380 and 3350 discs will be used but eventually only 3380 discs will prevail. This change will have an impact on data storage at London and users are advised to reassess their data storage habits.

i) U.K. CCP5 participants are reminded that at Daresbury Laboratory the S.E.R.C. has available an FPS 164 attached processor, currently with 3 MAX (matrix accelerator) boards, which is available for grant supported computing within the Science Board Community. U.K. research groups are invited to apply informally to Dr. J.E. Inglesfield, at Daresbury, for time on the FPS for benchmarking purposes. If the FPS proves viable for a given project, a formal application may be made. Advice on using the FPS at Daresbury may be obtained from the User Interface Group (in the person of Dr. D. Taylor) or from the CCP5 representative, Dr. W. Smith. A document entitled 'Using the FPS Attached Processor at the Daresbury Laboratory' by Dr. M.F. Guest is available from the User Interface Group.

j) Concurrent Computing Initiative at Daresbury.

A new advanced research computing group has been set up at Daresbury to develop large scale computational codes for a number of scientific applications in a highly parallel computer environment. Central to this project is the recently acquired FPS-T20 hypercube, 16 powerful processors interconnected as a 4 dimensional hypercube or tesseract

(hence T for tesseract and $20_8=2^4$).

Each node of the T-machine comprises an Inmos transputer, 1Mb of memory and 3 Weitek floating point chips, giving a realisable 12 Mflop performance per node on highly vectorised code. Thus the total peak performance is 192 Mflop with 160 Mflop having already been demonstrated on a distributed matrix multiply. The machine is frontended by a Micro-Vax II running Micro-VMS 4.4 and initially the T-machine is only programmable in occam. However a software release scheduled for August will provide both Fortran and C, and also move to a Unix frontend operating system.

Plans for upgrades are already in place and it is anticipated that the system will evolve to be an extremely powerful Fortran engine for both scalar and vector code. It is only by exploiting the inherent parallelism of most scientific applications that one can hope for the several orders of magnitude increase in computational power required by many applications. Current supercomputers with small numbers of expensive processors sharing memory cannot realise the 1000 fold increase in power offered by highly concurrent, distributed memory architectures.

R.J.Harrison and M.F.Guest.

SERC, Daresbury Laboratory, Warrington WA4 4AD.

k) The international conference entitled 'The Impact of Supercomputers on Chemistry' is due to take place in the University of London from 13 - 16 April 1987. The conference will cover all aspects of the use of supercomputers; applications and methodology. Some of the talks that are to be given include: "Supercomputer Simulation of Solids" by M.J. Gillan, "The Impact of Quantum Monte Carlo on Chemistry" by D.M. Ceperley, "Quantum Chemistry, Statistical Mechanics and Fluid Dynamics" by E. Clementi, "Chemistry as a Many-Particle Problem" by H.J.C. Berendsen and "Protein Structure, Computer Simulation and Graphics in the Design of Novel Molecules" by T.L. Blundell.

Enquiries regarding the conference should be addressed to Dr. J. Altmann, ISOC 87 Secretariat, Room 209, University of London Computer Centre, 20 Guildford Street, London WC1N 1DZ. (Telex 8953011).

l) Readers may also be interested in the Seventh Summer School on Computing Techniques in Physics, which takes place

on 9th-18th. June 1987 at Bechyne Castle near Tabor, Czechoslovakia. This year's theme is 'Microcomputers in Physics'. The organiser is Dr. J. Nadrchal, address; Summer Schools on Computational Physics, Institute of Physics, Czechoslovak Academy of Sciences, Na Slovance 2, CS-180 40 Praha 8, Czechoslovakia.

m) The Institute of Physics is organising the 1987 Solid State Physics Conference for 16-18 December at the H.H. Wills Physics Laboratory, University of Bristol and it will contain a number of symposia that will undoubtedly be of interest to participants in CCP5. Among the intended symposia are: Statistical Mechanics of Disordered and Inhomogeneous Systems; Physics of Sputtering Processes; Polymers and Liquid Crystals; Scattering of Neutrons from Disordered Systems and The Physics of Solids at High Pressures. People interested in attending or contributing a paper should contact Dr. B.L. Gyorffy or Dr. R. Evans at the H.H. Wills Physics Laboratory, University of Bristol, Tyndall Avenue, Bristol BS8 1TL. We will provide more information in our next newsletter.

n) It has been suggested to the CCP5 Editors that it would be useful for CCP5 participants to have a list of the EARN and BITNET addresses of other members of the project. For this purpose we would be happy to receive from our readers their addresses, which we shall collect here at Daresbury for publication in the newsletter at a later date. So that we don't fall foul of the 'Data Protection Act', it should be understood by all participants that the contributed addresses will be available to everyone.

o) The CCP5 Program Library. Documents and programs are available free of charge to academic centres upon application to Dr. W. Smith, TCS Division, S.E.R.C. Daresbury Laboratory, Daresbury, Warrington WA4 4AD, U.K. Daresbury Laboratory. Listings of programs are available if required but it is recommended that magnetic tapes (to be supplied by the applicant) be used. It may also be possible to transfer a small number of programs over the JANET network to other computer centres in the U.K.. Please note that use of inappropriate packing for magnetic tapes (e.g. padded bags) may result in them being considered unusable by Daresbury Computing Division and returned without the required software. Please ensure that these forms of packaging are not used. A list of programs available is presented in the following pages.

We should also like to remind our readers that we would welcome contributions to the Program Library. The Library exists to provide support for the research efforts of everyone active in computer simulation and to this end we are always pleased to extend the range of software available. If any of our readers have any programs they would like to make available, please would they contact Dr. Smith.

THE CCP5 PROGRAM LIBRARY

MDATOM by S. M. Thompson. M.D. simulation of monatomic fluids.

HMDIAT by S. M. Thompson. M.D. simulation of homonuclear diatomic fluids.

MDLIN by S. M. Thompson. M.D. simulation of linear molecules.

MDLINQ by S. M. Thompson. M.D. simulation of linear molecules with quadrupole.

MDTETRA by S. M. Thompson. M.D. simulation of tetrahedral molecules.

MDPOLY by S. M. Thompson. M.D. simulation of polyatomic molecule fluids.

ADMIXT by W. Smith. M.D. simulation of monatomic molecule mixtures.

MDNACL by W. Smith. M.D. simulation of alkali halide salts.

MDMIXT by W. Smith. M.D. simulation of polyatomic molecule mixtures.

MDMULP by W. Smith. M.D. simulation of polyatomic molecule mixtures.

MDMPOL by W. Smith & D. Fincham. M.D. simulation of polyatomic molecule mixtures.

MDZOID by W. Smith & K. Singer. M.D. simulation of ellipsoidal gaussian molecules.

DENCOR by W. Smith. Calculation of density correlation functions.

CURDEN by W. Smith. Calculation of current density correlation functions.

HLJ1 by D. M. Heyes. M.D. simulation of atomic fluids.

HLJ2 by D. M. Heyes. M.D. simulation of atomic fluids (with VACF).

HLJ3 by D. M. Heyes. M.D. simulation of atomic fluids (link-cells).

HLJ4 by D. M. Heyes. M.D. simulation of atomic fluids (constant T or P).

HLJ5 by D. M. Heyes. M.D. simulation of atomic fluids (shifted force).

HLJ6 by D. M. Heyes. M.D. simulation of atomic fluids (Toxvaerd algorithm).

MCRPM by D. M. Heyes. M.C. simulation of electrolytes (restricted primitive model).

SURF by D. M. Heyes. M.D. simulation of model alkali halide lamina.

HSTOCH by W. F. van Gunsteren & D. M. Heyes. MD or SD simulation of molecules in vacuo or rectangular cell with solvent or lattice.

MDATOM by D. Fincham. M.D. simulation of atomic fluids.

MDDIAT by D. Fincham. M.D. simulation of diatomic molecule fluids.

MDDIATQ by D. Fincham. M.D. simulation of diatomic fluids.
 MDIONS by D. Fincham & N. Anastasiou. M.D. simulation of electrolytes.
 MDMANY by D. Fincham & W. Smith. Simulation of rigid polyatomic molecules with/without fractional charges.
 CARLOS by B. Jonsson & S. Romano. M.C. simulation of a polyatomic solute molecule in an aqueous cluster.
 MCN by N. Corbin. M.C. simulation of atomic fluids (Metropolis).
 SCN by N. Corbin. M.C. simulation of atomic fluids (Rossky, Friedman and Doll).

STATIC SIMULATION CODES

CASCADE by M. Leslie and W. Smith. Calculates the structure and energy of a defect in an ionic crystal for a given potential model.
 THBREL Determines the minimum energy configuration of a perfect lattice for a given potential.
 THBFIT Empirically fits a potential to experimentally observed properties of a perfect lattice.
 SYMLAT Determines the minimum energy configuration of a perfect lattice for a given potential.
 THBPHON Calculates phonon dispersion curves for ionic crystals with three-body terms in the potential.

University of London

ROYAL HOLLOWAY AND BEDFORD NEW COLLEGE

Egham Hill, Egham, Surrey TW20 0EX

Telephone Egham (0784) 34455

ext. 3401

DEPARTMENT OF CHEMISTRY
The Bourne Laboratory

Dr. D.M. Heyes

10th March 1987

Dear Colleague,

At the 5th March 1987 CCP5 executive committee meeting it was decided to look into the possibility of holding a CCP5 sponsored workshop on "Transport Processes".

I am writing to you to discover if you would be interested in attending such a meeting and if so the directions you would like it to take. An obvious area would be equilibrium and non-equilibrium MD routes to viscosity, self-diffusion and thermal conductivity of fluids. Transport in solids and algorithms for performing the molecular dynamics are other areas of interest. We would appreciate your opinions on this matter, and also on the months in this year most suitable to you.

Best wishes,

Yours sincerely,

David Heyes

Telex : 935504 Fax : 0784-37520

MEETING ANNOUNCEMENT:

THERMODYNAMICS AND STATISTICAL MECHANICS OF SMALL SYSTEMS

A meeting is being organised by the Royal Society of Chemistry Statistical Mechanics and Thermodynamics Group at the University of York on September 24-25 1987.

The topics for discussion include small system behaviour in porous media and as crystallites, clusters and micelles. Both theoretical and experimental aspects will be discussed and contributions are invited.

The speakers include Dr. R.J. Evans (Bristol), Prof. D.H. Everett (Bristol) and Prof. J.S. Rowlinson (Oxford).

Please contact Dr. J.H.R. Clarke, Department of Chemistry, U.M.I.S.T., Sackville Street, Manchester M60 1QD as soon as possible for details.

P.G. Francis

REVIEW OF THE CCP5 MEETING ON 'NEW HORIZONS IN THE COMPUTER SIMULATION OF CONDENSED PHASES' HELD AT UMIST 8 - 9 JANUARY 1987.

W. Smith

As the title of the meeting suggests, the latest CCP5 meeting was a forward looking affair, which was intended to highlight areas of computer simulation which seem likely to open new avenues of enquiry. The meeting was organised by the departing CCP5 Chairman, Dr. J.H.R. Clarke, with the assistance of his colleagues at the Department of Chemistry and was attended by seventy simulation scientists from the United Kingdom, Europe and North America. As is always the case with CCP5 meetings, many interesting questions were raised and Dr. Clarke wisely allowed sufficient time for discussion. As a result it was generally held to be a stimulating and informative meeting.

The first contributor to the meeting was Prof. W.L. Jorgensen (Purdue University), who spoke on the subject of computer simulations of organic reactions. The objective of this work was the determination of the rates of organic reactions from first principles. Prof. Jorgensen first considered the S_N2 reaction beloved of organic chemists, in particular, the reaction between chloromethane CH_3Cl and chloride anion Cl^- . S_N2 reactions are bimolecular and stereospecific, they do not have intermediates and have an activation energy that increases with the polarity of the solvent. Prof. Jorgensen's approach involved the calculation of the gas phase potential reaction surface followed by its modification to include solute-solvent effects and finally a simulation of the reaction along a suitable reaction coordinate. The first step was achieved through ab initio methods; calculating the energy of the transition state (i.e. the $[Cl \cdots CH_3 \cdots Cl]^-$ complex) over the reaction coordinate. The solute-solvent interactions were determined by the ab initio calculation of clusters of solvent molecules with the reactive species. The purpose of the simulation (by Monte Carlo method) was to determine the density function $g(r_c)$ giving the probability of the reacting species being at a separation r_c on the reaction coordinate. From this the 'reaction force' $w(r_c) = -kT \ln(g(r_c))$ could be obtained. The simulation involved the reacting system plus 250 'TIP4P' water molecules within a periodic system. The MC procedure required the simulation to be done for several overlapping 'windows' of the reaction coordinate, using a biasing technique to sample each window adequately. The final set of overlapping windows gave the full $g(r_c)$ function. The NTP ensemble required 1-2M

equilibration steps and 2-6M steps averaging for each window. In the simulation the potential functions used were Lennard-Jones 12-6 functions with fractional charges to represent the Coulombic forces. However the parameters were described as functions of the reaction coordinate, fitted to the ab initio calculations.

The results for the example reaction showed that the relatively complicated energy barrier of the gas phase reaction was replaced by a single, higher peak in aqueous solution. These results tally with experimental observations, which show the gas reaction to be much faster, and that the solvated reaction has no intermediates. Interestingly, in the less polar solvent DMF, the energy barrier showed more structure, particularly the occurrence of minima suggestive of intermediates. (This latter suggestion has some experimental NMR evidence in its favour.) Calculations of the energy and number distributions showed that the number of solvating water molecules about the reacting species is constant throughout the reaction, contrary to expectations. The relative reduction of the solvation energy of the transition state (a factor contributing to the energy barrier increase) is due to a weakening of the solute-solvent bonding, not to the expected reduction in solvent coordination numbers. Further calculations using thermodynamic perturbation theory indicate that the structure of the transition state in gas and solvent is very similar.

Prof. Jorgensen concluded with a study of the S_N1 reaction of t-butyl chloride with Cl^- and showed that the formation of a contact ion pair $[t-Bu]^+Cl^-$ was clearly apparent, but the evidence for a solvent separated ion pair was much less clear. Throughout his talk Prof. Jorgensen presented excellent graphics revealing the stereochemical aspects of the reactions being described.

Dr. Sv. Toxvaerd (Panum Institute) described his studies in the implementation of the Generalised Langevin Equation (GLE). He was concerned to show how the GLE could be used to simulate the system of a polymer in solution for the fraction of the cost of a full molecular dynamics simulation but with comparable accuracy. The system he considered was a nonmer consisting of Lennard-Jones atoms bound by harmonic forces and suspended in a solution with 509 Lennard-Jones atoms. Following the work of Bossis, Quentrec and Boon, the GLE for this system is written as:

$$m \frac{dv_i(t)}{dt} = -m \int dt' M_i(t-t') v_i(t') + r_i + e_i(t)$$

where v_i is the velocity of the i 'th. bead, r_i the stochastic force, e_i the bead-bead forces on bead i , and $M_i(t)$ the memory function. The division of the forces into stochastic and bead-bead contributions is only valid if the correlation function $\langle v_i \cdot r_i \rangle$ is zero throughout. Dr. Toxvaerd has checked this using simulation and has established that it is true except at short times, where small deviations occur.

Once the memory functions M_i are known, the GLE can be solved numerically and Dr. Toxvaerd described his procedure for obtaining them. The memory function for the whole system is replaced by a Mori expansion. These terms cannot in general be resolved into intramolecular and solvent contributions, with the exception of the first term K_1 . Nevertheless the higher terms can be calculated by molecular dynamics via the time derivatives of the velocities. Explicit consideration of the first two terms of the Mori expansion; comparison of the values for the pure solvent with values obtained for the polymer solution allows physically reasonable approximations to the memory functions M_i to be made. The memory functions are then written as:

$$M_i(t) = K_{1,i}(s,t) \exp(-K_2 t^2/2)$$

In which $K_{1,i}(s,t)$ is the instantaneous value of the solvent dependent part of K_1 (which is dependent on the bead number i) and K_2 is the second Mori coefficient for the pure solvent. This approximation has been implemented in a simulation program using the Fourier method of Wang and Uhlenbeck and the results have been compared with molecular dynamics and a simpler Langevin Equation (LE) simulation.

The results show that small differences arise between the MD and the GLE and LE simulations of the bead-bead bondlength distribution, but the bead velocity autocorrelation functions are well reproduced. The approximate GLE method was shown to give essentially exact results for the centre of mass diffusion. Studies of the end-to-end distance R^2 and the radius of gyration S^2 and the corresponding correlation functions showed that very long runs are necessary for good statistics (prohibitively for the MD simulations) and about 10^6 timesteps are required for the molecule to establish the mean shape. Within the limitations on the studies occasioned by these difficulties, the MD and stochastic dynamics results were good. A study of the effects of the solvent on the molecular conformation was undertaken, by means of a comparison with a 'free molecule' simulation of the monomer. For the state points studied, the molecule was rather globular, and no great effects of the

solvent were seen. It was concluded that the GLE method described was a very useful method for studying polymers in solution.

Dr. D. Rapaport (Bar-Ilan University) spoke on the new subject of microscale hydrodynamics, by means of which he has successfully simulated some of the macroscopic features of fluid flow. He began by describing the kinds of phenomena seen in real flow experiments, such as fluid flow past a stationary object. At low flow rates (low Reynolds number) the flow is laminar (Stokes flow). As the flow increases stationary eddies form. At yet higher rates the eddies become unstable and move away from the object (vortex shedding) and give rise to a wake pattern known as the vortex or von Karman street. At very high flow rates turbulence results.

Dr. Rapaport described how he had simulated many of these features in a molecular dynamics simulation on the parallel assembly of FPS attached processors at IBM, Kingston, New York. The simulation was of two-dimensional flow past a fixed cylinder of 25 nm diameter. The fluid consisted of 170,000 'soft sphere' molecules in a 150 nm wide system with periodic boundaries. A fluid density of 0.83 and a flow speed of 330 m/s was used. A force field of 10^{11} g was employed to drive the flow and rethermalisation of the molecules at the down-field boundary was used to maintain the periodic conditions. The entire simulation cell was partitioned into 'slabs' to permit efficient parallelisation of the system and reduce data communication between processors to a minimum. The resulting flow pattern was presented graphically. The simulation cell was partitioned into a 60x60 grid and a short time average of the particle velocities in each sub-cell was calculated at intervals of 5 timesteps. The average velocity vector for each sub-cell was then plotted as a field of flow. Approximately 100 hours of CPU time were required on a system of four FPS 264 processors in parallel.

The results showed clearly the development of the different regimes of flow behaviour described above i.e. laminar flow, eddy formation, incipient instability of the eddies and vortex shedding and finally the von Karman street. The method clearly has great potential in elucidating the microscopic features that underlie the continuum behaviour of fluids in flow. Dr. Rapaport concluded by discussing the limitations of the method (short times, small spatial scales and subsonic velocities) and possible extensions of the method to three dimensional simulations.

Dr. M.P. Allen (University of Bristol) described his work done in a collaboration with D. Frenkel (Utrecht) on the molecular dynamics simulation of the isotropic to nematic phase transition in a system of hard ellipsoids. (The phase diagram for this system had been obtained previously by Frenkel and Mulder.) Dr Allen described the technique by which the hard ellipsoid system was simulated. (A detailed description of this appears in CCP5 Info. Quart. No. 23 p.28.) In practice the system simulated consisted of 144 ellipsoids and generated 10^6 collisions per hour on a CYBER 205 computer. The ellipsoids studied had axial ratios (a/b) of 2 and 3. The objective was to calculate single and collective time-dependent orientational correlation functions in the region of the isotropic-nematic phase transition and deduce the values of parameters important in the DeGennes and Keyes/Kivelson theories.

The single particle correlation functions were defined as:

$$c_L^S(t) = \langle P_L(\hat{e}_i(0) \cdot \hat{e}_i(t)) \rangle$$

and the collective correlations functions as:

$$c_L^C(t) = \sum_j \langle P_L(\hat{e}_i(0) \cdot \hat{e}_j(t)) \rangle / \sum_j \langle P_L(\hat{e}_i(0) \cdot \hat{e}_j(0)) \rangle$$

in which P_L is a Legendre polynomial of order L and \hat{e}_i is a unit vector aligned along the major axis of the ellipsoid. These functions were calculated for $L=1$ and 2 at various densities. The functions showed exponential decay at long times in the expected manner. The characteristic decay times τ for the collective correlation functions were significantly longer than the decay times for the single particle correlations, particularly near nematic-isotropic phase transition. The behaviour of the ratio τ_2^S/τ_2^C was shown in a plot against density and revealed a critical slowing down effect at the transition for the axial ratio 3. This slowing down was interpreted using the Kivelson expression:

$$\tau_2^S/\tau_2^C = (1+j_2)/(1+g_2)$$

where g_2 is the static second-rank Kirkwood factor and j_2 is a dynamical quantity that is describable in terms of the orientational memory functions. Calculations of the factor $(1+g_2)^{-1}$ for different densities allowed the above equation to be used to calculate j_2 , which was found to be insensitive to the density. The calculated mean values were given as $1+j_2=0.74\pm 0.16$ (a/b=2) and $1+j_2=0.92\pm 0.05$ (a/b=3). These results find experimental evidence in their

favour from the work of Flygare et al. on p - methoxy - benzylidene - n - butylaniline (MBBA), which provided an estimate of $l+j_2$ of the order 0.7 and also supported the mean field predictions of DeGennes that g_2 diverges at the critical density. Further systematic investigations of these and related systems are planned.

Dr. D. Tildesley (University of Southampton) described the method he has devised with his collaborators M.R. Stapleton (Southampton) and N. Quirke (BP Sunbury-on-Thames) for simulating variable polydispersed liquids which he defined as systems in which the dispersed particles (micelles) change size when the density of the system changes. The systems examined were the hard-sphere system and a modified Lennard-Jones system, in both of which every particle possessed a different effective diameter (i.e. $\sigma_{eff} = (\sigma_i + \sigma_j)/2$). The simulation method was a modified Monte Carlo algorithm in which the sizes of the particles as well as their locations were changed as the simulation proceeded. The sampling of the particle sizes at each MC move was taken from an appropriate particle size distribution function, which in this case was the Shultz function:

$$f_z(\sigma) = z!^{-1}((z+1)/\sigma)^{z+1} \sigma^z \exp(-((z+1)/\sigma)\sigma) \text{ if } [a < \sigma < b]$$

$$= 0 \text{ otherwise}$$

Where z is a parameter controlling the form of the function. (The distribution function changes with density and temperature for the LJ fluid.) A move was accepted or rejected according to the change in the inter-micellar interaction energy resulting from the move. Special care was needed for the long range corrections. Applications of this method to hard spheres showed that as the number density of the particles increases, they tend to become smaller, while for Lennard-Jones particles there is an initial average increase in the particle diameters followed by a decrease at higher densities. The initial increase was thought to be due to the attractive component of the LJ potential, which encouraged particles to approach, but this effect was outweighed by repulsive component when the density of the particles was high.

For fluids in which the distribution function is symmetric an adequate description is possible using a monodispersed fluid with the appropriate packing fraction, which is a function of the density and the third moment of the equilibrium distribution of diameters; M_3 . A mean field theory for M_3 has been developed in collaboration with Dr. T. Sluckin (Southampton). The chemical potential

distributions of the hard sphere polydispersed fluids were calculated by means of Widom's particle insertion method with lattices of different particle sizes. The scaled particle theory of Salacuse and Stell gave good agreement with the simulated chemical potential obtained with both low density distribution functions.

The second day began with a talk by Dr. D. Price on the computer simulation of minerals important in the understanding of geological processes. The minerals of greatest significance being silicates. It is extremely important that the properties of these materials be fully understood before realistic models of geological processes such as plate tectonics can be developed. The problem however is that these materials most commonly exist in regions of high pressure and temperature in the earth's core and is not practicable to perform suitable experiments in the laboratory. (Perovskite, a high pressure phase of magnesium silicate, for example is thought to constitute 40 percent of the earth by volume, but it is still very little understood). Computer simulation is therefore an essential tool for investigating such materials. The simulations need to provide information on several properties of the minerals to assist the geologist; the equation of state, (P vs. T) phase diagram, bulk and shear moduli, rheological properties and diffusion processes among them.

A number of methods are available for modelling these systems: (i) the ab initio Hartree-Fock method employing periodic boundaries (accurate but limited to small systems); (ii) ab initio calculations on SiO clusters to obtain the potential energy surfaces (good for some properties, though there is a loss in accuracy resulting from harmonic fitting of the surfaces); (iii) modified electron gas potentials with fitting of the potential surface using empirical (e.g. Buckingham) potentials (poor estimates of the elastic constants result) and (iv) The ionic model using fitted empirical potentials, plus coulombic interactions and the shell model (massless charged shell coupled to ion core by harmonic spring) to account for polarisability. The whole periodic crystal structure is modelled. To offset the effects of covalency non integer charges and bond bending forces are sometimes used.

Of the four methods described only the last is presently useful for providing geological data. The method is employed in the large lattice simulation programs CASCADE, THREL and THBPHON and give excellent crystal structure data and elasticity constants (e.g. the predicted structures of olivine, spinel and ilmenite are very good).

Phonon data are also available and allow the calculation of heat capacity, Gruneisen parameters, thermal expansion, entropy and defect entropies. Also the phonon spectrum has been found to be in close agreement with infra red and Raman frequencies. The eigenvectors provide details of the vibrational motions of the ions which is not easily available from experiments. The calculations of the entropies and heat capacities have been found to be very accurate. Thermal expansion coefficients and the phase diagrams of minerals have been obtained by this method and compare fairly well with experimental data. Lastly, defect energy calculations, using the CASCADE program, have allowed reasonable estimates of the diffusion rates of ions.

In the future Dr. Price expects to see the ionic model being used successfully to determine unknown phases and the incorporation of better three-body forces such as the Axilrod-Teller type. It is also expected that ab initio and modified electron gas potentials will be sufficiently advanced to assist with this type of work.

Dr. R. Lynden-Bell (University of Cambridge) delivered a talk on the subject of phase transitions and the fluctuations in orientationally disordered crystals. Dr. Lynden-Bell has used molecular dynamics to study these phenomena in a variety of systems; CH_4 , CBr_4 , CCl_4 , Na^+CN^- , tBuBr and SF_6 . The value of the molecular dynamics approach is that it complements existing experimental and theoretical methods and naturally introduces the effects of temperature and entropy. The theory of ordering studied by Dr. Lynden-Bell is the Landau theory, in which the free energy of a system is given by the expression:

$$G = G_0 + a(T-T_c)Y^2 + cY^4$$

In which Y is the 'order parameter'. This function has a single minimum for temperatures $T > T_c$ (indicating a stable, ordered phase) a very broad minimum when $T = T_c$ (indicating incipient instability) and a double minimum when $T < T_c$ (indicating a new preferential ordering). The order parameter Y is itself a function of the orientation of the molecules and is an irreducible representation of the space group of the crystal. Its spatial Fourier transform is represented by:

$$Y(\underline{k}) = \sum_i^{\text{mol}} \exp(i\underline{k} \cdot \underline{r}_i) K_{\Gamma M}^L(\Omega_i)$$

In which the function $K_{\Gamma M}^L(\Omega_i)$ is a combination of spherical harmonics (e.g. cubic harmonics) appropriate to the crystal symmetry. The index L indicates the type of ordering; thus

L=1 implies dipolar ordering, L=2 quadrupolar ordering and L=3 octupolar ordering. Fluctuations in the order parameter Y (i.e. $\langle Y^2 \rangle$) are largest when a change in phase is incipient and plots of $\langle Y^2 \rangle$ vs. k vector provide information on the crystallographic features of the transition e.g. a peak at $k=0$ (Γ point) indicates no change in the number of molecules in the unit cell, at k -zone boundary implies a doubling (or similar) in the unit cell and at intermediate k values, an incommensurate phase is expected.

Dr. Lynden-Bell illustrated these principles with examples of the orientational properties of sodium nitrite and nitrogen on graphite. The concluding example was an investigation of tertiary butyl bromide (tBuBr) to determine whether the ordering was dipolar (L=1) or octupolar (L=3); the former indicating the dominance of the molecular dipole in the crystal structure and the latter the dominance of the tetrahedral geometry of the molecule. A comparison of the $\langle Y^2 \rangle$ vs. k plots of these two alternatives with the corresponding plots of CBr_4 , showed the closer resemblance to the octupolar case.

Prof. C.R.A. Catlow (University of Keele / Daresbury) described the investigation of the structure and conductivity of β alumina. Two forms of β alumina exist; the stoichiometric form $\text{Na}_2\text{O} \cdot 11\text{Al}_2\text{O}_3$ and the non-stoichiometric form $(\text{Na}_2\text{O}) \cdot 1.311\text{Al}_2\text{O}_3$. The believed structures are complex spinel blocks in which the sodium atoms occur in layers at alternate hexagonal sites (the so-called Beavers-Ross (BR) sites). The sodium ions are mobile in these layers and give rise to conductivity. The non-stoichiometric sodium ions reside in bridging locations between oxygen ions. The structures were investigated both by static lattice methods (which used a shell model for the ions) and by molecular dynamics (which used a rigid ion model and the Ewald sum to handle the long range forces). The objectives were to determine the role of non-stoichiometry in conduction, the migration mechanism, the structure of the conduction plane and the temperature dependence of the conduction.

The static lattice calculations were performed with the CASCADE program and showed that the interstitial sodium ions occupied the 'anti-Beavers-Ross' sites and that the conduction path for the sodium ions was different in the non-stoichiometric and stoichiometric cases. The molecular dynamics were performed with the program FUNGUS, which was adapted for the parallel assembly of FPS processors at Kingston N.Y.. The simulations used 500 ions per MD cell; the stoichiometric case possessing 24 Na^+ ions. The non-stoichiometric cases had 4 and 8 more Na^+ ions (plus

compensating O^{2-} ions) respectively.

The stoichiometric cases showed the sodium ions to be at the BR sites but NO diffusion (i.e. conductivity) was evident in the temperature region 600K to 900K, though the Na^+ ions showed large thermal parameters. The non-stoichiometric ($+4Na^+$) case showed signs of diffusion at 600K and definite diffusion at 900K. The interstitial sodium ions occupied anti-BR sites with some displacement due to strong interaction with the excess O^{2-} ions. There was some evidence of conduction paths. Surprisingly, the non-stoichiometric ($+8Na^+$) case showed significantly less diffusion and extensive interactions between the excess sodium and oxygen ions. The excess sodium ions were displaced towards the mid O-O bridge.

The conclusions of this work were that the stoichiometric β alumina is not a sodium ion conductor, while in the highly non-stoichiometric form conduction is suppressed. In the low non-stoichiometric form the excess sodium ions occupy anti-BR sites, but are displaced towards the oxygens at higher non-stoichiometry. Conduction by the sodium ions appears to involve distinct pathways.

The session on the quantum mechanical methods of simulation began with a talk by Dr. M. Gillan (Harwell) on the Path Integral Monte Carlo (PIMC) method, which is a method for simulating quantum systems at nonzero temperatures and has its origins in the path integral formulation of quantum mechanics due to Feynman. Beginning with the one dimensional, single particle case, Dr. Gillan explained the derivation of the method. At the core of the method is the mathematical function known as the density matrix, in terms of which the partition function of a system and the mean value of any static (equilibrium) observable quantity may be calculated. The form of the density matrix is:

$$\rho(x, x'; \beta) = \sum_n \varphi_n^*(x) \exp(-\beta E_n) \varphi_n(x')$$

which closely resembles the Green's function or propagator used in the path integral formulation of Feynman. This means that a substitution of the time variable by $i\beta\hbar/2\pi$ permits the formal use of the path integral method to calculate the partition function for the system of interest. In reality the density function has only been derived for two cases; the free particle and the harmonic potential, and no analytical solution is known for complex potentials or many-body systems. However Feynman's formalism allows the partitioning of a given path integral into shorter paths

over the span of which the harmonic potential may be used as a good approximation. (This is the so-called short-time approximation.) If this strategy is adopted, it emerges that the resultant expression for the partition function is formally equivalent to that of a classical model ring polymer in which the atoms are coupled by harmonic springs with a force constant of $2mP(\pi)^2/(\hbar\beta)^2$ and each atom is in a potential field of strength V/P , where P is the number of atoms in the ring and V is the potential in the original Hamiltonian.

The physical significance of the polymer is that collectively it represents the delocalisation of the parent particle in agreement with the Uncertainty Principle. The effective reduction of the potential energy by the factor $1/P$ is what enables the parent particle to tunnel into classically forbidden regions. The spring constant is largest when the temperature or the mass of the parent particle is large. This means that the ring polymer reduces in extent (i.e. reduced delocalisation) under these circumstances, which comply with the classical limit. Generalisation of the method to many dimensions and many particles is easily accomplished (but may be computationally expensive).

Simulations using this method require only that some realistic way of generating the configurations of the isomorphic polymer be available. Either molecular dynamics, Monte Carlo or stochastic dynamics methods may be employed. (No significance should be attached to the time dependent aspects of MD and SD.) Dr. Gillan prefers the stochastic method over MD since it avoids the problems of ergodicity encountered in MD when the spring constants are large. Dr. Gillan gave a number of examples of the application of the method and concluded with his own work on the dissolution of hydrogen in metals in which he has been able to reproduce many of the observed features of these remarkable systems.

Dr. B.H. Wells (University of Oxford) continued the quantum mechanical theme in a talk on the new developments that have taken place in the Green's Function Monte Carlo method. Dr. Wells began by describing the two basic methods currently in use; the Variational Monte Carlo and Diffusion Monte Carlo methods. Both are methods for dealing with the electronic wavefunctions of molecular systems. The Variational method is used to calculate the energy of a system from a known wavefunction. In the method an initial set of electron positions are moved using a gaussian propagator and the moves accepted according to a probability based on the square of the wavefunction. The so-called local

energy $E_L(r')$ is calculated and accumulated at each step, following an equilibration period. In Diffusion Monte Carlo the wavefunction itself is calculated, starting from an initial trial wavefunction. In this method the gaussian propagation is augmented by an exponential branching term, for which a trial state energy is required and is updated as the simulation proceeds. Excited states can be calculated provided the initial wavefunction possesses the correct symmetry. The standard method uses the 'fixed node' approximation, which requires the nodes of the wavefunction to be accurately known. This knowledge prevents negative configuration densities arising in the sampling procedure. Configurations which cross a node boundary in the course of the calculation are deleted, but provided the trial energy is close to the actual state energy, the population of configurations will be maintained. The method has been successfully applied to calculations of the correlation energy in small molecules.

Among the advances in the methods that have been taking place recently are the removal of the short-time approximation in the Green's function propagator, the removal of the 'fixed-node' constraint and the development of the differential Monte Carlo method to enable calculations of physical quantities, such as the molecular dipole, which depend on small energy differences. (The essential feature of the latter method is the calculation of two distinct systems, one the minor perturbation of the other, using the same random walk to reduce the variances of the calculated property.) Example applications of the methods described by Dr. Wells included the $e_2^+e^-$ bound system, e^-H^+ systems at a dielectric surface (for both of which good binding energies were obtained) and the calculations of H_3^+ and H_5^+ systems with and without the Born-Oppenheimer approximation.

The final speaker of the meeting was Dr. D. Fincham (University of York) who provided an overview of the current developments in parallel computing. He began by defining the concept of a parallel computer, which is a computer possessing more than one processing unit, all or some of which could be active simultaneously under the direction of the user's program. Parallel computers may be SIMD (Single Instruction Multiple Datastream) or MIMD (Multiple Instruction Multiple Datastream) machines.

There are currently three main architectures for parallel computers. The first of these is the mainframe which has a small number of processors with shared memory (e.g. Cray XMP/48). The second is the switching network, in

which there are separate processors and interconnected memories (e.g. ULTRA). In the third architecture, the processors are connected but have their own memory boards. An important feature of all parallel computers is the topology of the connections between processing units. Dr. Fincham gave four examples; the square array (e.g. ICL DAP), the hypercube (e.g. INTEL hypercube), the pyramid (devised by the University of Erlangen) and the Waterloo Ring. Each of these topologies has its own features and advantages e.g. the hypercube, which possesses 2^n nodes with n connections to each, has other connectivities as subsets of this, and these can be exploited in specific applications such as the periodic boundary.

A recent significant development in the parallel computing concept has been the introduction of the TRANSPUTER, which is a high performance microprocessor which has been specifically designed to be linked to other TRANSPUTERS so creating a parallel processor. Each processor is itself a complete computer, which will soon acquire a hardware floating point unit (currently floating point arithmetic is handled by software). By way of an example application Dr. Fincham described a molecular dynamics application of the TRANSPUTER, that was suitable for simulating biological molecules. Other recent significant developments are: the INTEL iPSC and iPSC/VX computers which have the hypercube topology and up to 128 processors; the FPS T series, which has hypercube topology and vector processors at each node; the ICL Mini-DAP, a SIMD machine, and the Reconfigurable Processor Array (RPA) at Southampton, in which the processing elements can be dynamically configured to suit the application.

Dr. Fincham concluded with a few comments on the (unlikely) transportability of parallel codes and the commercial pressures likely to encourage the development of larger machines with fewer processors instead of smaller multi-processor machines.

Throughout the meeting a number of poster displays were on view. These included:

"Dynamical Simulation of a Model Reverse Micelle", by D. Brown and J.H.R. Clarke.

"Simulated Annealing of Atomic Clusters", by L.T. Wille.

"Computer Simulation of Plasticity in Minerals", by M. Heggie, R. Jones and Y. Zhang.

"Simulation Studies on Adsorption of Methanol over Silicalite", by R. Vetrivel, C.R.A. Catlow and E.A. Colbourn.

"Computer Simulation of Zeolite Stability and Structure", by R.A. Jackson.

"Generalised Hydrodynamics and Eigenmodes for a Lennard-Jones Fluid", by C. Bruin.

"Computer Simulation of a Classical Fluid with Internal Quantum States", by J. Talbot, P. Ballone, Ph. de Smedt and J.L. Lebowitz.

Acknowledgements

I am grateful to Dr Julian Clarke, Professor Konrad Singer and to several of the speakers for their assistance in preparing this review. Any errors that are present are entirely of my own making.

Computer Simulation Studies of Zeolite Structure and Stability

R. A. Jackson and C. R. A. Catlow

Department of Chemistry, University of Keele, Keele, Staffs. ST5 5BG.

Introduction

Atomistic computer simulation techniques have been applied to a range of zeolites, with a view to predicting the structures and relative stabilities of these materials.

Details of the Method

The atomistic computer simulation technique that has been used in this work is lattice energy minimisation. This technique enables lattice energies to be calculated and the associated minimum energy structures to be obtained. It relies on the availability of interatomic potentials to describe the interactions of atoms in the crystal; the potential model used for zeolites has been applied to a wide range of ionic and semi-ionic solids, and the parameters have been used in a successful simulation of the structures of α -quartz and its polymorphs⁽¹⁾. An important feature of this model is the inclusion of a 3 body term to account for covalent effects in silicon-oxygen bonding.

A second derivative minimisation method has been used in all calculations; it should be noted that in the case of silicalite, the CRAY X-MP/48 was required for the calculations.

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on Molecular Dynamics, Monte Carlo and Lattice Simulations of Condensed Phases.

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THE UNIVERSITY OF CHICAGO

ANNUAL REPORT OF THE BOARD OF TRUSTEES

1958

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Editorial.

It is probably part of the natural cycle of things that it occasionally becomes necessary to appeal for contributions for the CCP5 Newsletter. Such is the case now. This does not however indicate that there is a diminishing interest in the newsletter itself, as rarely a day goes by when we are not asked to add more names to our mailing list. It is therefore important to remind our readers that the newsletter is very dependent on their enthusiasm if it is to keep going.

The CCP5 newsletter is an interesting publication in a number of ways: it published regularly, articles appear quickly (within three months or less), it has an international readership and it is informal enough to allow the publication of novel or speculative ideas. Added to these is the fact that there are surprisingly few publications that discuss the computational details of simulation methods warts and all. For these reasons it should be well worth supporting.

So to encourage our readership we announce that we would be pleased to receive contributions from anyone active in the field of computer simulation of condensed phases, anywhere in the world! We would also be particularly pleased to publish material from postgraduate researchers, who often have some original ideas to contribute to the computational art. We await your contributions!

In view of this appeal it is a special pleasure to thank the contributors to this month's issue. Long may they prosper!

Contributors:

D. Brown	Department of Chemistry, U.M.I.S.T., Sackville St., Manchester M60 1QD.
C.R.A. Catlow R.A Jackson	Department of Chemistry, University of Keele, Keele, Staffs. ST5 5BC.
J.E. Quinn	DAP Support Unit, Queen Mary College, Mile End Road, London E1 4NS.
W. Smith	Theory and Computational Science Division, S.E.R.C. Daresbury

Laboratory, Daresbury, Warrington
WA4 4AD.

L. Woodcock

Department of Chemical Engineering,
The University of Bradford,
Bradford, Yorkshire BD7 1DP.

General News.

a) CCP5 is to organise a meeting entitled "Industrial Applications of Computer Simulation". The meeting will take place at the Royal Institution in London from 11th. to 13th. January 1988. The purpose of the meeting is to discuss the application of the simulation methods associated with CCP5 (i.e. molecular dynamics, Monte Carlo, energy minimisation etc.) to industrially significant problems. The organisers of the meeting will be Prof. C.R.A. Catlow and Dr. N. Quirke. Readers interested in taking part in this meeting should write to: Dr. N. Quirke, B.P. Research Centre, Chertsey Road, Sunbury-on-Thames, Middlesex TW16 7LN.

Further news of this meeting will be announced later.

b) CCP5 is also to organise a meeting on the subject of bio-organic applications of computer simulation, which provisionally, is scheduled for Easter 1988. More details of this will be revealed as the organisation develops.

c) At the last meeting of the CCP5 Steering Committee which took place on 9th. January the composition of the Executive Committee was changed in accordance with the usual procedure. Professor C.R.A. Catlow was elected as the new CCP5 Chairman to replace Dr. J.H.R. Clarke. In addition Drs. R. Lynden-Bell, S. Parker and M.P. Allen were elected to the Executive Committee. The outgoing members being Drs. N. Quirke and P.A. Madden. In addition to these changes; Dr. M. Leslie has taken the responsibility of Secretary to CCP5 and Dr. W. Smith will now handle the CCP5 Program Library. The composition of the Executive Committee is thus:

Professor C.R.A. Catlow (Keele University and Daresbury)-
Chairman,

Dr. M.P. Allen (Bristol University)

Dr. D.M. Heyes (Royal Holloway and Bedford New College)

Dr. M. Leslie (Daresbury) - Secretary

Dr. R. Lynden-Bell (Cambridge University)

Dr. S. Parker (Bath University)

Dr. W. Smith (Daresbury) - Program Librarian

This would seem to be the place to thank the past members of the Executive Committee for their work on behalf of the CCP5 community. In particular it is appropriate to thank Julian Clarke, under whose chairmanship for the past three years CCP5 has been extremely successful.

d) The CCP5 Executive Committee wish to announce to all U.K. participants in CCP5 that it has set aside a part of the CCP5 budget for the purpose of assisting collaborative work in the U.K. simulation community. These funds are to assist with expenses for travel between collaborating centres. So far, CCP5 has agreed to assist the following collaborating groups:

J. Clarke and D. Tildesley (Transport Properties),

S. Parker and G. Price (Lattice Simulations),

M. Allen and G. Evans (Hard Ellipsoids),

K. Singer and W. Smith (Quantum Simulations).

Other groups who wish to be considered for this assistance should write to the CCP5 Chairman, Prof. C.R.A. Catlow, Department of Chemistry, University of Keele, Keele, Staffs. ST5 5BG.

e) The CCP5 Executive Committee also wishes to remind U.K. participants that it has an annual allocation of Cray time at each of the centres: London (Cray 1s) and Rutherford (Cray XMP-48), which is available for the development of simulation programs prior to a grant allocation. At present CCP5 is allocated 15 hours a year at London and 5 hours quarterly at Rutherford. Readers who wish to use some of this allocation should write to the CCP5 Secretary, Dr. M. Leslie, TCS Division, SERC Daresbury Laboratory, Daresbury, Warrington WA4 4AD.

f) The Rutherford and Appleton Computer Centre Cray XMP/48 is now up and running. It is currently accessible to a few privileged users, who are running jobs to test the machine in a working environment before it becomes available to the wider academic community in April.

In time it is expected that the CCP5 Program Library will be modified to run on the XMP and users who would like to see any specific programs converted immediately should contact the CCP5 Program Librarian, Dr. W. Smith.

g) The University of Manchester Regional Computing Centre has described the new VM/CMS service which they will be running soon on the Amdahl 5890/300. It was originally intended to run the three domains of the Multiple Domain Facility to deal with: (i) the service currently supplied by the Amdahl 470 V/8; (ii) a national VM/CMS service and (iii) a Manchester local service. However only two domains are possible with the current hardware and so the last two services will run initially on one domain. Some disruption

of the existing service is expected 20-23 March to reconnect the peripherals of the 470 V/8 to the 5890/300, followed by acceptance tests. The experimental VM/CMS service is expected to be available in early May. A useful description of the VM/CMS service appears in the March UMRCC Newsletter.

An important aspect of the operation of all the national computing centres (London, Manchester and Rutherford-Appleton) will be the new Peer Review system for the allocation of computing resources. Everyone involved in scientific computing in the UK should become familiar with this. The UMRCC newsletter for March 1987 and the ULCC newsletter for February 1987 both carry details of the agreed procedure and are thus recommended reading.

h) The University of London Computing Centre has now upgraded its front end computer to an Amdahl 5890. The second phase of this upgrade will involve the introduction of new disc storage devices (Amdahl 6380 disc drives, which are 3380 compatible). During this phase both 3380 and 3350 discs will be used but eventually only 3380 discs will prevail. This change will have an impact on data storage at London and users are advised to reassess their data storage habits.

i) U.K. CCP5 participants are reminded that at Daresbury Laboratory the S.E.R.C. has available an FPS 164 attached processor, currently with 3 MAX (matrix accelerator) boards, which is available for grant supported computing within the Science Board Community. U.K. research groups are invited to apply informally to Dr. J.E. Inglesfield, at Daresbury, for time on the FPS for benchmarking purposes. If the FPS proves viable for a given project, a formal application may be made. Advice on using the FPS at Daresbury may be obtained from the User Interface Group (in the person of Dr. D. Taylor) or from the CCP5 representative Dr. W. Smith. A document entitled 'Using the FPS Attached Processor at the Daresbury Laboratory' by Dr. M.F. Guest is available from the User Interface Group.

j) Concurrent Computing Initiative at Daresbury.

A new advanced research computing group has been set up at Daresbury to develop large scale computational codes for a number of scientific applications in a highly parallel computer environment. Central to this project is the recently acquired FPS-T20 hypercube, 16 powerful processors interconnected as a 4 dimensional hypercube or tesseract

(hence T for tesseract and $20_8=2^4$).

Each node of the T-machine comprises an Inmos transputer, 1Mb of memory and 3 Weitek floating point chips, giving a realisable 12 Mflop performance per node on highly vectorised code. Thus the total peak performance is 192 Mflop with 160 Mflop having already been demonstrated on a distributed matrix multiply. The machine is frontended by a Micro-Vax II running Micro-VMS 4.4 and initially the T-machine is only programmable in occam. However a software release scheduled for August will provide both Fortran and C, and also move to a Unix frontend operating system.

Plans for upgrades are already in place and it is anticipated that the system will evolve to be an extremely powerful Fortran engine for both scalar and vector code. It is only by exploiting the inherent parallelism of most scientific applications that one can hope for the several orders of magnitude increase in computational power required by many applications. Current supercomputers with small numbers of expensive processors sharing memory cannot realise the 1000 fold increase in power offered by highly concurrent, distributed memory architectures.

R.J.Harrison and M.F.Guest.

SERC, Daresbury Laboratory, Warrington WA4 4AD.

k) The international conference entitled 'The Impact of Supercomputers on Chemistry' is due to take place in the University of London from 13 - 16 April 1987. The conference will cover all aspects of the use of supercomputers; applications and methodology. Some of the talks that are to be given include: "Supercomputer Simulation of Solids" by M.J. Gillan, "The Impact of Quantum Monte Carlo on Chemistry" by D.M. Ceperley, "Quantum Chemistry, Statistical Mechanics and Fluid Dynamics" by E. Clementi, "Chemistry as a Many-Particle Problem" by H.J.C. Berendsen and "Protein Structure, Computer Simulation and Graphics in the Design of Novel Molecules" by T.L. Blundell.

Enquiries regarding the conference should be addressed to Dr. J. Altmann, ISOC 87 Secretariat, Room 209, University of London Computer Centre, 20 Guildford Street, London WC1N 1DZ. (Telex 8953011).

l) Readers may also be interested in the Seventh Summer School on Computing Techniques in Physics, which takes place

on 9th-18th June 1987 at Bechyne Castle near Tabor, Czechoslovakia. This year's theme is 'Microcomputers in Physics'. The organiser is Dr. J. Nadrchal, address; Summer Schools on Computational Physics, Institute of Physics, Czechoslovak Academy of Sciences, Na Slovance 2, CS-180 40 Praha 8, Czechoslovakia.

m) The Institute of Physics is organising the 1987 Solid State Physics Conference for 16-18 December at the H.H. Wills Physics Laboratory, University of Bristol and it will contain a number of symposia that will undoubtedly be of interest to participants in CCP5. Among the intended symposia are: Statistical Mechanics of Disordered and Inhomogeneous Systems; Physics of Sputtering Processes; Polymers and Liquid Crystals; Scattering of Neutrons from Disordered Systems and The Physics of Solids at High Pressures. People interested in attending or contributing a paper should contact Dr. B.L. Gyorffy or Dr. R. Rvans at the H.H. Wills Physics Laboratory, University of Bristol, Tyndall Avenue, Bristol BS8 1TL. We will provide more information in our next newsletter.

n) It has been suggested to the CCP5 Editors that it would be useful for CCP5 participants to have a list of the EARN and BITNET addresses of other members of the project. For this purpose we would be happy to receive from our readers their addresses, which we shall collect here at Daresbury for publication in the newsletter at a later date. So that we don't fall foul of the 'Data Protection Act', it should be understood by all participants that the contributed addresses will be available to everyone.

o) The CCP5 Program Library. Documents and programs are available free of charge to academic centres upon application to Dr. W. Smith, TCS Division, S.E.R.C. Daresbury Laboratory, Daresbury, Warrington WA4 4AD, U.K. Daresbury Laboratory. Listings of programs are available if required but it is recommended that magnetic tapes (to be supplied by the applicant) be used. It may also be possible to transfer a small number of programs over the JANET network to other computer centres in the U.K.. Please note that use of inappropriate packing for magnetic tapes (e.g. padded bags) may result in them being considered unusable by Daresbury Computing Division and returned without the required software. Please ensure that these forms of packaging are not used. A list of programs available is presented in the following pages.

We should also like to remind our readers that we would welcome contributions to the Program Library. The Library exists to provide support for the research efforts of everyone active in computer simulation and to this end we are always pleased to extend the range of software available. If any of our readers have any programs they would like to make available, please would they contact Dr. Smith.

THE CCP5 PROGRAM LIBRARY

MDATOM by S. M. Thompson. M.D. simulation of monatomic fluids.

HMDIAT by S. M. Thompson. M.D. simulation of homonuclear diatomic fluids.

MDLIN by S. M. Thompson. M.D. simulation of linear molecules.

MDLIQ by S. M. Thompson. M.D. simulation of linear molecules with quadrupole.

MDTETRA by S. M. Thompson. M.D. simulation of tetrahedral molecules.

MDPOLY by S. M. Thompson. M.D. simulation of polyatomic molecule fluids.

ADMIXT by W. Smith. M.D. simulation of monatomic molecule mixtures.

MDNAEL by W. Smith. M.D. simulation of alkali halide salts.

MDMIXT by W. Smith. M.D. simulation of polyatomic molecule mixtures.

MDMULP by W. Smith. M.D. simulation of polyatomic molecule mixtures.

MDMPOL by W. Smith & D. Fincham. M.D. simulation of polyatomic molecule mixtures.

MDZOID by W. Smith & K. Singer. M.D. simulation of ellipsoidal gaussian molecules.

DENCOR by W. Smith. Calculation of density correlation functions.

CURDEN by W. Smith. Calculation of current density correlation functions.

HLJ1 by D. M. Heyes. M.D. simulation of atomic fluids.

HLJ2 by D. M. Heyes. M.D. simulation of atomic fluids (with VACF).

HLJ3 by D. M. Heyes. M.D. simulation of atomic fluids (link-cells).

HLJ4 by D. M. Heyes. M.D. simulation of atomic fluids (constant T or P).

HLJ5 by D. M. Heyes. M.D. simulation of atomic fluids (shifted force).

HLJ6 by D. M. Heyes. M.D. simulation of atomic fluids (Toxvaerd algorithm).

MCRPM by D. M. Heyes. M.C. simulation of electrolytes (restricted primitive model).

SURF by D. M. Heyes. M.D. simulation of model alkali halide lamina.

HSTOCH by W. F. van Gunsteren & D. M. Heyes. MD or SD simulation of molecules in vacuo or rectangular cell with solvent or lattice.

MDATOM by D. Fincham. M.D. simulation of atomic fluids.

MDDIAT by D. Fincham. M.D. simulation of diatomic molecule fluids.

MDDIATQ by D. Fincham. M.D. simulation of diatomic fluids.
 MDIONS by D. Fincham & N. Anastasiou. M.D. simulation of electrolytes.
 MDMANY by D. Fincham & W. Smith. Simulation of rigid polyatomic molecules with/without fractional charges.
 CARLOS by B. Jonsson & S. Romano. M.C. simulation of a polyatomic solute molecule in an aqueous cluster.
 MCN by N. Corbin. M.C. simulation of atomic fluids (Metropolis).
 SCN by N. Corbin. M.C. simulation of atomic fluids (Rossky, Friedman and Doll).

STATIC SIMULATION CODES

CASCADE by M. Leslie and W. Smith. Calculates the structure and energy of a defect in an ionic crystal for a given potential model.
 THBREL Determines the minimum energy configuration of a perfect lattice for a given potential.
 THBFIT Empirically fits a potential to experimentally observed properties of a perfect lattice.
 SYMLAT Determines the minimum energy configuration of a perfect lattice for a given potential.
 THBPHON Calculates phonon dispersion curves for ionic crystals with three-body terms in the potential.

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Telephone Egham (0784) 34455

ext. 3401

DEPARTMENT OF CHEMISTRY
The Bourne Laboratory

Dr. D.W. Heyes

10th March 1987

Dear Colleague,

At the 5th March 1987 CCP5 executive committee meeting it was decided to look into the possibility of holding a CCP5 sponsored workshop on "Transport Processes".

I am writing to you to discover if you would be interested in attending such a meeting and if so the directions you would like it to take. An obvious area would be equilibrium and non-equilibrium MD routes to viscosity, self-diffusion and thermal conductivity of fluids. Transport in solids and algorithms for performing the molecular dynamics are other areas of interest. We would appreciate your opinions on this matter, and also on the months in this year most suitable to you.

Best wishes,

Yours sincerely,

D.W. Heyes

Telex : 935504 Fax : 0784-37520

MEETING ANNOUNCEMENT:

THERMODYNAMICS AND STATISTICAL MECHANICS OF SMALL SYSTEMS

A meeting is being organised by the Royal Society of Chemistry Statistical Mechanics and Thermodynamics Group at the University of York on September 24-25 1987.

The topics for discussion include small system behaviour in porous media and as crystallites, clusters and micelles. Both theoretical and experimental aspects will be discussed and contributions are invited.

The speakers include Dr. R.J. Evans (Bristol), Prof. D.H. Everett (Bristol) and Prof. J.S. Rowlinson (Oxford).

Please contact Dr. J.H.R. Clarke, Department of Chemistry, U.M.I.S.T., Sackville Street, Manchester M60 1QD as soon as possible for details.

P.G. Francis

REVIEW OF THE CCP5 MEETING ON 'NEW HORIZONS IN THE COMPUTER SIMULATION OF CONDENSED PHASES' HELD AT UMIST 8 - 9 JANUARY 1987.

W. Smith

As the title of the meeting suggests, the latest CCP5 meeting was a forward looking affair, which was intended to highlight areas of computer simulation which seem likely to open new avenues of enquiry. The meeting was organised by the departing CCP5 Chairman, Dr. J.H.R. Clarke, with the assistance of his colleagues at the Department of Chemistry and was attended by seventy simulation scientists from the United Kingdom, Europe and North America. As is always the case with CCP5 meetings, many interesting questions were raised and Dr. Clarke wisely allowed sufficient time for discussion. As a result it was generally held to be a stimulating and informative meeting.

The first contributor to the meeting was Prof. W.L. Jorgensen (Purdue University), who spoke on the subject of computer simulations of organic reactions. The objective of this work was the determination of the rates of organic reactions from first principles. Prof. Jorgensen first considered the S_N2 reaction beloved of organic chemists, in particular, the reaction between chloromethane CH_3Cl and chloride anion Cl^- . S_N2 reactions are bimolecular and stereospecific, they do not have intermediates and have an activation energy that increases with the polarity of the solvent. Prof. Jorgensen's approach involved the calculation of the gas phase potential reaction surface followed by its modification to include solute-solvent effects and finally a simulation of the reaction along a suitable reaction coordinate. The first step was achieved through ab initio methods; calculating the energy of the transition state (i.e. the $[Cl \cdots CH_3 \cdots Cl]^-$ complex) over the reaction coordinate. The solute-solvent interactions were determined by the ab initio calculation of clusters of solvent molecules with the reactive species. The purpose of the simulation (by Monte Carlo method) was to determine the density function $g(r_c)$ giving the probability of the reacting species being at a separation r_c on the reaction coordinate. From this the 'reaction force' $w(r_c) = -kT \ln(g(r_c))$ could be obtained. The simulation involved the reacting system plus 250 'TIP4P' water molecules within a periodic system. The MC procedure required the simulation to be done for several overlapping 'windows' of the reaction coordinate, using a biasing technique to sample each window adequately. The final set of overlapping windows gave the full $g(r_c)$ function. The NTP ensemble required 1-2M

equilibration steps and 2-6M steps averaging for each window. In the simulation the potential functions used were Lennard-Jones 12-6 functions with fractional charges to represent the Coulombic forces. However the parameters were described as functions of the reaction coordinate, fitted to the ab initio calculations.

The results for the example reaction showed that the relatively complicated energy barrier of the gas phase reaction was replaced by a single, higher peak in aqueous solution. These results tally with experimental observations, which show the gas reaction to be much faster, and that the solvated reaction has no intermediates. Interestingly, in the less polar solvent DMF, the energy barrier showed more structure, particularly the occurrence of minima suggestive of intermediates. (This latter suggestion has some experimental NMR evidence in its favour.) Calculations of the energy and number distributions showed that the number of solvating water molecules about the reacting species is constant throughout the reaction, contrary to expectations. The relative reduction of the solvation energy of the transition state (a factor contributing to the energy barrier increase) is due to a weakening of the solute-solvent bonding, not to the expected reduction in solvent coordination numbers. Further calculations using thermodynamic perturbation theory indicate that the structure of the transition state in gas and solvent is very similar.

Prof. Jorgensen concluded with a study of the S_N1 reaction of t-butyl chloride with Cl^- and showed that the formation of a contact ion pair $[t-Bu]^+ \cdot Cl^-$ was clearly apparent, but the evidence for a solvent separated ion pair was much less clear. Throughout his talk Prof. Jorgensen presented excellent graphics revealing the stereochemical aspects of the reactions being described.

Dr. S. Toxvaerd (Panum Institute) described his studies in the implementation of the Generalised Langevin Equation (GLE). He was concerned to show how the GLE could be used to simulate the system of a polymer in solution for the fraction of the cost of a full molecular dynamics simulation but with comparable accuracy. The system he considered was a nonmer consisting of Lennard-Jones atoms bound by harmonic forces and suspended in a solution with 509 Lennard-Jones atoms. Following the work of Bossis, Quentrec and Boon, the GLE for this system is written as:

$$m \frac{dv_i(t)}{dt} = - \int dt' M_i(t-t') v_i(t') + r_i + e_i(t)$$

where \underline{v}_i is the velocity of the i 'th. bead, \underline{r}_i the stochastic force, \underline{g}_i the bead-bead forces on bead i , and $M_i(t)$ the memory function. The division of the forces into stochastic and bead-bead contributions is only valid if the correlation function $\langle \underline{v}_i \cdot \underline{r}_i \rangle$ is zero throughout. Dr. Toxvaerd has checked this using simulation and has established that it is true except at short times, where small deviations occur.

Once the memory functions M_i are known, the GLE can be solved numerically and Dr. Toxvaerd described his procedure for obtaining them. The memory function for the whole system is replaced by a Mori expansion. These terms cannot in general be resolved into intramolecular and solvent contributions, with the exception of the first term K_1 . Nevertheless the higher terms can be calculated by molecular dynamics via the time derivatives of the velocities. Explicit consideration of the first two terms of the Mori expansion; comparison of the values for the pure solvent with values obtained for the polymer solution allows physically reasonable approximations to the memory functions M_i to be made. The memory functions are then written as:

$$M_i(t) = K_{1,i}(s,t) \exp(-K_2 t^2/2)$$

In which $K_{1,i}(s,t)$ is the instantaneous value of the solvent dependent part of K_1 (which is dependent on the bead number i) and K_2 is the second Mori coefficient for the pure solvent. This approximation has been implemented in a simulation program using the Fourier method of Wang and Uhlenbeck and the results have been compared with molecular dynamics and a simpler Langevin Equation (LE) simulation.

The results show that small differences arise between the MD and the GLE and LE simulations of the bead-bead bondlength distribution, but the bead velocity autocorrelation functions are well reproduced. The approximate GLE method was shown to give essentially exact results for the centre of mass diffusion. Studies of the end-to-end distance R^2 and the radius of gyration S^2 and the corresponding correlation functions showed that very long runs are necessary for good statistics (prohibitively for the MD simulations) and about 10^6 timesteps are required for the molecule to establish the mean shape. Within the limitations on the studies occasioned by these difficulties, the MD and stochastic dynamics results were good. A study of the effects of the solvent on the molecular conformation was undertaken, by means of a comparison with a 'free molecule' simulation of the monomer. For the state points studied, the molecule was rather globular, and no great effects of the

solvent were seen. It was concluded that the GLE method described was a very useful method for studying polymers in solution.

Dr. D. Rapaport (Bar-Ilan University) spoke on the new subject of microscale hydrodynamics, by means of which he has successfully simulated some of the macroscopic features of fluid flow. He began by describing the kinds of phenomena seen in real flow experiments, such as fluid flow past a stationary object. At low flow rates (low Reynolds number) the flow is laminar (Stokes flow). As the flow increases stationary eddies form. At yet higher rates the eddies become unstable and move away from the object (vortex shedding) and give rise to a wake pattern known as the vortex or von Karman street. At very high flow rates turbulence results.

Dr. Rapaport described how he had simulated many of these features in a molecular dynamics simulation on the parallel assembly of FPS attached processors at IBM, Kingston, New York. The simulation was of two-dimensional flow past a fixed cylinder of 25 nm diameter. The fluid consisted of 170,000 'soft sphere' molecules in a 150 nm wide system with periodic boundaries. A fluid density of 0.83 and a flow speed of 330 m/s was used. A force field of 10^{11} g was employed to drive the flow and rethermalisation of the molecules at the down-field boundary was used to maintain the periodic conditions. The entire simulation cell was partitioned into 'slabs' to permit efficient parallelisation of the system and reduce data communication between processors to a minimum. The resulting flow pattern was presented graphically. The simulation cell was partitioned into a 60x60 grid and a short time average of the particle velocities in each sub-cell was calculated at intervals of 5 timesteps. The average velocity vector for each sub-cell was then plotted as a field of flow. Approximately 100 hours of CPU time were required on a system of four FPS 264 processors in parallel.

The results showed clearly the development of the different regimes of flow behaviour described above i.e. laminar flow, eddy formation, incipient instability of the eddies and vortex shedding and finally the von Karman street. The method clearly has great potential in elucidating the microscopic features that underlie the continuum behaviour of fluids in flow. Dr. Rapaport concluded by discussing the limitations of the method (short times, small spatial scales and subsonic velocities) and possible extensions of the method to three dimensional simulations.

Dr. M.P. Allen (University of Bristol) described his work done in a collaboration with D. Frenkel (Utrecht) on the molecular dynamics simulation of the isotropic to nematic phase transition in a system of hard ellipsoids. (The phase diagram for this system had been obtained previously by Frenkel and Mulder.) Dr Allen described the technique by which the hard ellipsoid system was simulated. (A detailed description of this appears in CCP5 Info. Quart. No. 23 p.28.) In practice the system simulated consisted of 144 ellipsoids and generated 10^6 collisions per hour on a CYBER 205 computer. The ellipsoids studied had axial ratios (a/b) of 2 and 3. The objective was to calculate single and collective time-dependent orientational correlation functions in the region of the isotropic-nematic phase transition and deduce the values of parameters important in the DeGennes and Keyes/Kivelson theories.

The single particle correlation functions were defined as:

$$c_L^S(t) = \langle P_L(e_i(0) \cdot e_i(t)) \rangle$$

and the collective correlations functions as:

$$c_L^C(t) = \sum_j \langle P_L(e_i(0) \cdot e_j(t)) \rangle / \sum_j \langle P_L(e_i(0) \cdot e_j(0)) \rangle$$

in which P_L is a Legendre polynomial of order L and e_i is a unit vector aligned along the major axis of the ellipsoid. These functions were calculated for L=1 and 2 at various densities. The functions showed exponential decay at long times in the expected manner. The characteristic decay times τ for the collective correlation functions were significantly longer than the decay times for the single particle correlations, particularly near nematic-isotropic phase transition. The behaviour of the ratio τ_2^S/τ_2^C was shown in a plot against density and revealed a critical slowing down effect at the transition for the axial ratio 3. This slowing down was interpreted using the Kivelson expression:

$$\tau_2^S/\tau_2^C = (1+j_2)/(1+g_2)$$

where g_2 is the static second-rank Kirkwood factor and j_2 is a dynamical quantity that is describable in terms of the orientational memory functions. Calculations of the factor $(1+g_2)^{-1}$ for different densities allowed the above equation to be used to calculate j_2 , which was found to be insensitive to the density. The calculated mean values were given as $1+j_2=0.74\pm 0.16$ (a/b=2) and $1+j_2=0.92\pm 0.05$ (a/b=3). These results find experimental evidence in their

favour from the work of Flygare et al. on p - methoxy - benzylidene - n - butylaniline (MBBA), which provided an estimate of $1+j_2$ of the order 0.7 and also supported the mean field predictions of DeGennes that g_2 diverges at the critical density. Further systematic investigations of these and related systems are planned.

Dr. D. Tildesley (University of Southampton) described the method he has devised with his collaborators M.R. Stapleton (Southampton) and N. Quirke (BP Sunbury-on-Thames) for simulating variable polydispersed liquids which he defined as systems in which the dispersed particles (micelles) change size when the density of the system changes. The systems examined were the hard-sphere system and a modified Lennard-Jones system, in both of which every particle possessed a different effective diameter (i.e. $\sigma_{eff} = (\sigma_i + \sigma_j)/2$). The simulation method was a modified Monte Carlo algorithm in which the sizes of the particles as well as their locations were changed as the simulation proceeded. The sampling of the particle sizes at each MC move was taken from an appropriate particle size distribution function, which in this case was the Shultz function:

$$f_z(\sigma) = z!^{-1}((z+1)/\sigma)^{z+1} \sigma^z \exp(-((z+1)/\sigma)\sigma) \text{ if } [a < \sigma < b]$$

$$= 0 \text{ otherwise}$$

Where z is a parameter controlling the form of the function. (The distribution function changes with density and temperature for the LJ fluid.) A move was accepted or rejected according to the change in the inter-micellar interaction energy resulting from the move. Special care was needed for the long range corrections. Applications of this method to hard spheres showed that as the number density of the particles increases, they tend to become smaller, while for Lennard-Jones particles there is an initial average increase in the particle diameters followed by a decrease at higher densities. The initial increase was thought to be due to the attractive component of the LJ potential, which encouraged particles to approach, but this effect was outweighed by repulsive component when the density of the particles was high.

For fluids in which the distribution function is symmetric an adequate description is possible using a monodispersed fluid with the appropriate packing fraction, which is a function of the density and the third moment of the equilibrium distribution of diameters; M_3 . A mean field theory for M_3 has been developed in collaboration with Dr. T. Sluckin (Southampton). The chemical potential

distributions of the hard sphere polydispersed fluids were calculated by means of Widom's particle insertion method with lattices of different particle sizes. The scaled particle theory of Salacuse and Stell gave good agreement with the simulated chemical potential obtained with both low density distribution functions.

The second day began with a talk by Dr. D. Price on the computer simulation of minerals important in the understanding of geological processes. The minerals of greatest significance being silicates. It is extremely important that the properties of these materials be fully understood before realistic models of geological processes such as plate tectonics can be developed. The problem however is that these materials most commonly exist in regions of high pressure and temperature in the earth's core and is not practicable to perform suitable experiments in the laboratory. (Perovskite, a high pressure phase of magnesium silicate, for example is thought to constitute 40 percent of the earth by volume, but it is still very little understood). Computer simulation is therefore an essential tool for investigating such materials. The simulations need to provide information on several properties of the minerals to assist the geologist; the equation of state, (P vs. T) phase diagram, bulk and shear moduli, rheological properties and diffusion processes among them.

A number of methods are available for modelling these systems: (i) the ab initio Hartree-Fock method employing periodic boundaries (accurate but limited to small systems); (ii) ab initio calculations on SiO₂ clusters to obtain the potential energy surfaces (good for some properties, though there is a loss in accuracy resulting from harmonic fitting of the surfaces); (iii) modified electron gas potentials with fitting of the potential surface using empirical (e.g. Buckingham) potentials (poor estimates of the elastic constants result) and (iv) The ionic model using fitted empirical potentials, plus coulombic interactions and the shell model (massless charged shell coupled to ion core by harmonic spring) to account for polarisability. The whole periodic crystal structure is modelled. To offset the effects of covalency non integer charges and bond bending forces are sometimes used.

Of the four methods described only the last is presently useful for providing geological data. The method is employed in the large lattice simulation programs CASCADE, THBREL and THBPHON and give excellent crystal structure data and elasticity constants (e.g. the predicted structures of olivine, spinel and ilmenite are very good).

Phonon data are also available and allow the calculation of heat capacity, Gruneisen parameters, thermal expansion, entropy and defect entropies. Also the phonon spectrum has been found to be in close agreement with infra red and Raman frequencies. The eigenvectors provide details of the vibrational motions of the ions which is not easily available from experiments. The calculations of the entropies and heat capacities have been found to be very accurate. Thermal expansion coefficients and the phase diagrams of minerals have been obtained by this method and compare fairly well with experimental data. Lastly, defect energy calculations, using the CASCADE program, have allowed reasonable estimates of the diffusion rates of ions.

In the future Dr. Price expects to see the ionic model being used successfully to determine unknown phases and the incorporation of better three-body forces such as the Axilrod-Teller type. It is also expected that ab initio and modified electron gas potentials will be sufficiently advanced to assist with this type of work.

Dr. R. Lynden-Bell (University of Cambridge) delivered a talk on the subject of phase transitions and the fluctuations in orientationally disordered crystals. Dr. Lynden-Bell has used molecular dynamics to study these phenomena in a variety of systems; CH_4 , CBr_4 , CCl_4 , $\text{Na}^+\cdot\text{CN}^-$, tBuBr and SF_6 . The value of the molecular dynamics approach is that it complements existing experimental and theoretical methods and naturally introduces the effects of temperature and entropy. The theory of ordering studied by Dr. Lynden-Bell is the Landau theory, in which the free energy of a system is given by the expression:

$$G = G_0 + a(T-T_c)Y^2 + cY^4$$

In which Y is the 'order parameter'. This function has a single minimum for temperatures $T > T_c$ (indicating a stable, ordered phase) a very broad minimum when $T = T_c$ (indicating incipient instability) and a double minimum when $T < T_c$ (indicating a new preferential ordering). The order parameter Y is itself a function of the orientation of the molecules and is an irreducible representation of the space group of the crystal. Its spatial Fourier transform is represented by:

$$Y(\underline{k}) = \sum_i^{\text{mol}} \exp(i\underline{k} \cdot \underline{r}_i) K_{\Gamma M}^L(\Omega_i)$$

In which the function $K_{\Gamma M}^L(\Omega_i)$ is a combination of spherical harmonics (e.g. cubic harmonics) appropriate to the crystal symmetry. The index L indicates the type of ordering; thus

$L=1$ implies dipolar ordering, $L=2$ quadrupolar ordering and $L=3$ octupolar ordering. Fluctuations in the order parameter Y (i.e. $\langle Y^2 \rangle$) are largest when a change in phase is incipient and plots of $\langle Y^2 \rangle$ vs. k vector provide information on the crystallographic features of the transition e.g. a peak at $k=0$ (Γ point) indicates no change in the number of molecules in the unit cell, at k -zone boundary implies a doubling (or similar) in the unit cell and at intermediate k values, an incommensurate phase is expected.

Dr. Lynden-Bell illustrated these principles with examples of the orientational properties of sodium nitrite and nitrogen on graphite. The concluding example was an investigation of tertiary butyl bromide (tBuBr) to determine whether the ordering was dipolar ($L=1$) or octupolar ($L=3$); the former indicating the dominance of the molecular dipole in the crystal structure and the latter the dominance of the tetrahedral geometry of the molecule. A comparison of the $\langle Y^2 \rangle$ vs. k plots of these two alternatives with the corresponding plots of CBr_4 , showed the closer resemblance to the octupolar case.

Prof. C.R.A. Catlow (University of Keele / Daresbury) described the investigation of the structure and conductivity of β alumina. Two forms of β alumina exist; the stoichiometric form $\text{Na}_2\text{O} \cdot 11\text{Al}_2\text{O}_3$ and the non-stoichiometric form $(\text{Na}_2\text{O}) \cdot 1.311\text{Al}_2\text{O}_3$. The believed structures are complex spinel blocks in which the sodium atoms occur in layers at alternate hexagonal sites (the so-called Beavers-Ross (BR) sites). The sodium ions are mobile in these layers and give rise to conductivity. The non-stoichiometric sodium ions reside in bridging locations between oxygen ions. The structures were investigated both by static lattice methods (which used a shell model for the ions) and by molecular dynamics (which used a rigid ion model and the Ewald sum to handle the long range forces). The objectives were to determine the role of non-stoichiometry in conduction, the migration mechanism, the structure of the conduction plane and the temperature dependence of the conduction.

The static lattice calculations were performed with the CASCADE program and showed that the interstitial sodium ions occupied the 'anti-Beavers-Ross' sites and that the conduction path for the sodium ions was different in the non-stoichiometric and stoichiometric cases. The molecular dynamics were performed with the program FUNGUS, which was adapted for the parallel assembly of FPS processors at Kingston N.Y.. The simulations used 500 ions per MD cell; the stoichiometric case possessing 24 Na^+ ions. The non-stoichiometric cases had 4 and 8 more Na^+ ions (plus

compensating O^{2-} ions) respectively.

The stoichiometric cases showed the sodium ions to be at the BR sites but NO diffusion (i.e. conductivity) was evident in the temperature region 600K to 900K, though the Na^+ ions showed large thermal parameters. The non-stoichiometric (+4 Na^+) case showed signs of diffusion at 600K and definite diffusion at 900K. The interstitial sodium ions occupied anti-BR sites with some displacement due to strong interaction with the excess O^{2-} ions. There was some evidence of conduction paths. Surprisingly, the non-stoichiometric (+8 Na^+) case showed significantly less diffusion and extensive interactions between the excess sodium and oxygen ions. The excess sodium ions were displaced towards the mid O-O bridge.

The conclusions of this work were that the stoichiometric β alumina is not a sodium ion conductor, while in the highly non-stoichiometric form conduction is suppressed. In the low non-stoichiometric form the excess sodium ions occupy anti-BR sites, but are displaced towards the oxygens at higher non-stoichiometry. Conduction by the sodium ions appears to involve distinct pathways.

The session on the quantum mechanical methods of simulation began with a talk by Dr. M. Gillan (Harwell) on the Path Integral Monte Carlo (PIMC) method, which is a method for simulating quantum systems at nonzero temperatures and has its origins in the path integral formulation of quantum mechanics due to Feynmann. Beginning with the one dimensional, single particle case, Dr. Gillan explained the derivation of the method. At the core of the method is the mathematical function known as the density matrix, in terms of which the partition function of a system and the mean value of any static (equilibrium) observable quantity may be calculated. The form of the density matrix is:

$$\rho(x, x'; \beta) = \sum_n \varphi_n^*(x) \exp(-\beta E_n) \varphi_n(x')$$

which closely resembles the Green's function or propagator used in the path integral formulation of Feynman. This means that a substitution of the time variable by $i\beta\hbar/2\pi$ permits the formal use of the path integral method to calculate the partition function for the system of interest. In reality the density function has only been derived for two cases; the free particle and the harmonic potential, and no analytical solution is known for complex potentials or many-body systems. However Feynman's formalism allows the partitioning of a given path integral into shorter paths

over the span of which the harmonic potential may be used as a good approximation. (This is the so-called short-time approximation.) If this strategy is adopted, it emerges that the resultant expression for the partition function is formally equivalent to that of a classical model ring polymer in which the atoms are coupled by harmonic springs with a force constant of $2mP(\pi)^2/(h\beta)^2$ and each atom is in a potential field of strength V/P , where P is the number of atoms in the ring and V is the potential in the original Hamiltonian.

The physical significance of the polymer is that collectively it represents the delocalisation of the parent particle in agreement with the Uncertainty Principle. The effective reduction of the potential energy by the factor $1/P$ is what enables the parent particle to tunnel into classically forbidden regions. The spring constant is largest when the temperature or the mass of the parent particle is large. This means that the ring polymer reduces in extent (i.e. reduced delocalisation) under these circumstances, which comply with the classical limit. Generalisation of the method to many dimensions and many particles is easily accomplished (but may be computationally expensive).

Simulations using this method require only that some realistic way of generating the configurations of the isomorphic polymer be available. Either molecular dynamics, Monte Carlo or stochastic dynamics methods may be employed. (No significance should be attached to the time dependent aspects of MD and SD.) Dr. Gillan prefers the stochastic method over MD since it avoids the problems of ergodicity encountered in MD when the spring constants are large. Dr. Gillan gave a number of examples of the application of the method and concluded with his own work on the dissolution of hydrogen in metals in which he has been able to reproduce many of the observed features of these remarkable systems.

Dr. B.H. Wells (University of Oxford) continued the quantum mechanical theme in a talk on the new developments that have taken place in the Green's Function Monte Carlo method. Dr. Wells began by describing the two basic methods currently in use; the Variational Monte Carlo and Diffusion Monte Carlo methods. Both are methods for dealing with the electronic wavefunctions of molecular systems. The Variational method is used to calculate the energy of a system from a known wavefunction. In the method an initial set of electron positions are moved using a gaussian propagator and the moves accepted according to a probability based on the square of the wavefunction. The so-called local

energy $E_L(r')$ is calculated and accumulated at each step, following an equilibration period. In Diffusion Monte Carlo the wavefunction itself is calculated, starting from an initial trial wavefunction. In this method the gaussian propagation is augmented by an exponential branching term, for which a trial state energy is required and is updated as the simulation proceeds. Excited states can be calculated provided the initial wavefunction possesses the correct symmetry. The standard method uses the 'fixed node' approximation, which requires the nodes of the wavefunction to be accurately known. This knowledge prevents negative configuration densities arising in the sampling procedure. Configurations which cross a node boundary in the course of the calculation are deleted, but provided the trial energy is close to the actual state energy, the population of configurations will be maintained. The method has been successfully applied to calculations of the correlation energy in small molecules.

Among the advances in the methods that have been taking place recently are the removal of the short-time approximation in the Green's function propagator, the removal of the 'fixed-node' constraint and the development of the differential Monte Carlo method to enable calculations of physical quantities, such as the molecular dipole, which depend on small energy differences. (The essential feature of the latter method is the calculation of two distinct systems, one the minor perturbation of the other, using the same random walk to reduce the variances of the calculated property.) Example applications of the methods described by Dr. Wells included the $e_2^+e^-$ bound system, e^-H^+ systems at a dielectric surface (for both of which good binding energies were obtained) and the calculations of H_3^+ and H_5^+ systems with and without the Born-Oppenheimer approximation.

The final speaker of the meeting was Dr. D. Fincham (University of York) who provided an overview of the current developments in parallel computing. He began by defining the concept of a parallel computer, which is a computer possessing more than one processing unit, all or some of which could be active simultaneously under the direction of the user's program. Parallel computers may be SIMD (Single Instruction Multiple Datastream) or MIMD (Multiple Instruction Multiple Datastream) machines.

There are currently three main architectures for parallel computers. The first of these is the mainframe which has a small number of processors with shared memory (e.g. Cray XMP/48). The second is the switching network, in

which there are separate processors and interconnected memories (e.g. ULTRA). In the third architecture, the processors are connected but have their own memory boards. An important feature of all parallel computers is the topology of the connections between processing units, Dr. Fincham gave four examples; the square array (e.g. ICL DAP), the hypercube (e.g. INTEL hypercube), the pyramid (devised by the University of Erlangen) and the Waterloo Ring. Each of these topologies has its own features and advantages e.g. the hypercube, which possesses 2^n nodes with n connections to each, has other connectivities as subsets of this, and these can be exploited in specific applications such as the periodic boundary.

A recent significant development in the parallel computing concept has been the introduction of the TRANSPUTER, which is a high performance microprocessor which has been specifically designed to be linked to other TRANSPUTERS so creating a parallel processor. Each processor is itself a complete computer, which will soon acquire a hardware floating point unit (currently floating point arithmetic is handled by software). By way of an example application Dr. Fincham described a molecular dynamics application of the TRANSPUTER, that was suitable for simulating biological molecules. Other recent significant developments are: the INTEL iPSC and iPSC/VX computers which have the hypercube topology and up to 128 processors; the FPS T series, which has hypercube topology and vector processors at each node; the ICL Mini-DAP, a SIMD machine, and the Reconfigurable Processor Array (RPA) at Southampton, in which the processing elements can be dynamically configured to suit the application.

Dr. Fincham concluded with a few comments on the (unlikely) transportability of parallel codes and the commercial pressures likely to encourage the development of larger machines with fewer processors instead of smaller multi-processor machines.

Throughout the meeting a number of poster displays were on view. These included:

"Dynamical Simulation of a Model Reverse Micelle", by D. Brown and J.H.R. Clarke.

"Simulated Annealing of Atomic Clusters", by L.T. Wille.

"Computer Simulation of Plasticity in Minerals", by M. Heggie, R. Jones and Y. Zhang.

"Simulation Studies on Adsorption of Methanol over Silicalite", by R. Vetrivel, C.R.A. Catlow and E.A. Colbourn.

- "Computer Simulation of Zeolite Stability and Structure", by R.A. Jackson.
- "Generalised Hydrodynamics and Eigenmodes for a Lennard-Jones Fluid", by C. Bruin.
- "Computer Simulation of a Classical Fluid with Internal Quantum States", by J. Talbot, P. Ballone, Ph. de Smedt and J.L. Lebowitz.

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Computer Simulation Studies of Zeolite Structure and Stability

R. A. Jackson and C. R. A. Catlow

Department of Chemistry, University of Keele, Keele, Staffs. ST5 5BG.

Introduction

Atomistic computer simulation techniques have been applied to a range of zeolites, with a view to predicting the structures and relative stabilities of these materials.

Details of the Method

The atomistic computer simulation technique that has been used in this work is lattice energy minimisation. This technique enables lattice energies to be calculated and the associated minimum energy structures to be obtained. It relies on the availability of interatomic potentials to describe the interactions of atoms in the crystal; the potential model used for zeolites has been applied to a wide range of ionic and semi-ionic solids, and the parameters have been used in a successful simulation of the structures of α -quartz and its polymorphs⁽¹⁾. An important feature of this model is the inclusion of a 3 body term to account for covalent effects in silicon-oxygen bonding.

A second derivative minimisation method has been used in all calculations; it should be noted that in the case of silicalite, the CRAY X-MP/48 was required for the calculations.

Results and Discussion

The work carried out falls into two sections. First, Na⁺ Zeolite A has been studied extensively, since reliable experimental data exist for this zeolite which enable the computer simulation methods to be tested. Secondly, the relative stabilities of a range of zeolites have been calculated.

In the case of Na⁺ zeolite A, a number of potential models have been used, ranging from a rigid ion description of all ions to the use of the shell model to describe oxygen ion polarisability. In the simulations it is most important that the positions of the framework ions be accurately reproduced; the non-framework cations (in this case Na⁺) are not so accurately located experimentally. This has been achieved in general, but some oxygen ion displacements are too large. Some refitting of the potential model (originally fitted to α -quartz) may be necessary.

In the case of calculations of relative stabilities, these have been obtained for the siliceous forms of zeolite A, faujasite and silicalite. Results are given in table 1. The corresponding value for α -quartz is given for comparison. It should be noted that these simulations were carried out to constant pressure, i.e. all strains, both internal and external, were removed.

An important extension of this work is the calculation of minimum Gibbs energies for these zeolites, as relative stability over a range of temperatures is important. This work is in progress.

Table 1

Lattice energies per SiO₂ unit after relaxation to constant pressure

	<u>lattice energy/eV</u>
Zeolite A	-123.66
Faujasite	-123.48
Silicalite	-123.89
quartz	-123.90

Acknowledgements

We are grateful to Shell Research BV for funding this work.

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SOME "AD HOC" NOTES ON ISOKINETIC AND (N,V,T) DYNAMICS

Les Woodcock

All equilibrium particle simulations comprise at least two stages; an equilibration period and a sample period. The relative computational requirements can vary enormously depending on the size of the system. For very large systems, $> \sim 10000$ particles the sample period for equilibrium properties may be instantaneous (one configuration). For very small systems eg. 2, any low energy configuration can be regarded as being in equilibrium. The computational overheads for equilibration of large systems with long relaxation time could be considerable; it is practical considerations such as these which will determine the choice of dynamical equations for the practitioner. The same applies to steady-state NEMD simulations.

When the non-equilibrium behaviour is of no interest the particles and momenta are set as close as possible to equilibrium. This usually involves placing the particles in a low-energy initial configuration and allocating a normal distribution of velocities. Irrespective of the mean and half-width of this initial velocity distribution $(\underline{v}_1 \dots \underline{v}_N)$ the requisite initial kinetic energy corresponding to an instantaneous temperature T_0 , defined by

$$3N^{\ddagger} k T_0 = \left(\sum_{i=1}^N \underline{p}_i^2 / m \right)_{t=0} = \langle \sum_{i=1}^N \underline{p}_i^2 / m \rangle \quad (1)$$

is exactly obtained by the scaling

$$(\underline{p}_1 \dots \underline{p}_N)_{T_0} = f_0 m (\underline{v}_1 \dots \underline{v}_N)_{\text{initial}} \quad (2)$$

The scaling factor f_0 is defined by

$$f_0^2 = \frac{3N^{\ddagger} k T_0}{\sum_{i=1}^N \underline{p}_i^2 / m} \quad (3)$$

The simplest method of isothermal constraint dynamics is to incorporate

this momentum renormalisation procedure continuously into the equations of motion. The objective is to generate isothermal (constant kinetic energy) dynamics with the minimum disturbance to the particle trajectories and momenta [1]. In the original continuous renormalisation method [2] a scaling factor $f(t)$ defined as above was incorporated directly into the second-order central difference algorithm of Verlet [3]. The finite difference equations (1) and (2) together with equation (7) of reference [2] are equivalent to the constraint equations of motion

$$d\underline{r}_i = \underline{p}_i dt \quad (4)$$

$$d\underline{p}_i = \underline{F}_i dt + (f(t) - 1)\underline{p}_i \quad (5)$$

$$f^2(t) = \frac{3N^{\ddagger} k T_0 m}{\sum_{i=1}^N (\underline{p}_i + \underline{F}_i dt)^2} \quad (6)$$

In equations (4) to (6) the force \underline{F}_i is $\nabla\phi(\underline{r}_i)$ where ϕ_i is the potential at \underline{r}_i .

If we substitute into equation (5) for f and square the bracket we obtain the continuous renormalisation equation for the acceleration

$$d\underline{p}_i = \underline{F}_i dt + \left(\left(\frac{3N^{\ddagger} k T_0 m}{\sum_{i=1}^N \underline{p}_i^2 + \sum_{i=1}^N 2\underline{p}_i \cdot \underline{F}_i dt + \sum_{i=1}^N (\underline{F}_i dt)^2} \right)^{1/2} - 1 \right) \underline{p}_i \quad (7)$$

substituting for T_0 from equation (1), expanding the inner bracket, and taking the limit $dt \rightarrow 0$ gives

$$d\underline{p}_i = \underline{F}_i dt - \frac{\sum_{i=1}^N \underline{p}_i \cdot \underline{F}_i}{\sum_{i=1}^N \underline{p}_i^2 / m} \underline{p}_i dt \quad (8)$$

Equation (8) is the isokinetic constraint equation of motion used by

Hoover et al. (4) and derived by Evans et al (5) from Gauss's principle of minimum constraint using a Lagrange undetermined multiplier. The derivation and some applications are reviewed by Evans and Morriss [1] where it is referred to as the method of Gaussian isokinetic dynamics.

Equations (1) to (8) show that the methods of continuous renormalisation [2] and Gaussian isokinetic dynamics are essentially the same. That the equations of motion (4-6) lead to isokinetic dynamics can be shown by deriving f from the constraint, additional to equation (1) for the average T_0 , that the time-derivative of the "instantaneous temperature" be zero, i.e.

$$\frac{dT(t)}{dt} = \frac{d}{dt} \frac{\sum_{i=1}^N p_i^2/m}{3Nk} = \frac{2 \sum_{i=1}^N p_i \cdot \dot{p}_i}{3Nkm} = 0 \quad (9)$$

substituting in equation (9) for $\dot{p}_i(t)$ from equation (5) gives the exact expression for f consistent with the constraint equation (9)

$$\frac{2}{3Nk} \sum_{i=1}^N \left(\frac{dp_i}{dt} + (f-1)p_i \right) p_i = 0 \quad (10)$$

$$\therefore \sum_{i=1}^N p_i \cdot \dot{p}_i = - \sum_{i=1}^N \frac{(f-1)p_i^2/m}{dt} \quad (11)$$

$$\therefore f = 1 + \frac{\sum_{i=1}^N \dot{p}_i \cdot p_i}{\sum_{i=1}^N p_i^2/m} \quad (12)$$

[Interpretation of f : time derivative of $(f-1)$ is a moderating factor which slows or speeds all the particles uniformly in momentum space].

The formal proof that this method generates truly isokinetic dynamics and, simultaneously reproduces the distribution in coordinate space of a canonical ensemble of statistical thermodynamics has been given recently by Nose [6].

Equation (6) is preferred to equation (12) for f in the implementation of these equations into a finite difference integration scheme because the continuous renormalisation has the advantages that:

- (i) the reference temperature T_0 enters the equations-of-motion
 - (ii) it is self-starting
 - (iii) it is more efficient for the equilibration period
- and (iv) it is self-stabilising up to much larger values of the finite time difference Δt .

That the continuous renormalisation method is practicable for the various finite difference algorithms normally employed for Newtonian MD has been demonstrated by many comparisons with N,V,E microcanonical calculations on similar systems. Actually it is generally found to stabilise the algorithms to larger values of the finite difference Δt . The net effect, however, of the temptation to use a larger Δt , is to unwittingly introduce some "Brownian motion" and this effect upon the time-dependent properties of the trajectories may be undesirable.

Several other schemes very closely related to the continuous renormalisation method have also appeared in the literature recently [7-10].

Brown and Clarke [7] have compared MD schemes, including a method which they refer to the "ad hoc" method of momentum scaling. Brown and Clarke's "ad hoc" method differs from the continuous renormalisation method in that the particles are slightly less constrained due to the absence of the $\underline{F}dt$ term in their scaling factor (equation (6)). The net effect of this is to introduce some fluctuations in T of the order $N^{-1/2}$, and linear in Δt ; this was evidenced by the comparisons of Brown and Clark with their Gaussian isokinetic calculations at the same mean T . It would appear that any "ad hoc" variation from equations (4-6) would to some degree loosen the constraint of constant kinetic energy, as also would too large a finite time increment.

Haille and Gupta [8] have introduced an alternative method of momentum scaling; they propose the equations of motion

$$\underline{dp}_i = \underline{F}_i dt \quad (13)$$

$$\underline{dr}_i = (\underline{p}_i/m) f_{HG} dt \quad (14)$$

where f_{HG} , the momentum scaling factor, is defined as f_o in equation (3). As with the "ad hoc" method of Brown and Clarke, the Haille-Gupta momentum scaling equations lead to deviations linear or first order, in dt , but they derive the equations from a generalised driving potential and indicate that the canonical coordinate distributions is generated. Nose [6], however, suggested that these H-G equations yield a modified canonical coordinate distribution with differences of the order $N^{-1/2}$.

Broughton et al. [9] have examined the effect of discontinuous renormalisation or periodic scaling.

Heyes et al [10] have added two further "ad hoc" methods to the growing list of variations of renormalisation techniques.

Any number of "ad hoc" methods can be derived but only where the algorithm reduces to equations (4-6) as $\Delta t \rightarrow 0$ does it correspond to the continuous renormalisation method of reference [2].

Berendsen et al [11] have derived a set of constraint equations of motion for isothermal (but not isokinetic) behaviour. They begin with the Langevin equation for macromolecule behaviour but remove the stochastic element. The Berendsen equations are isothermal, deterministic but not time reversible, and in the present notation take the form

$$\underline{dp}_i = \underline{F}_i dt + \gamma \left(\frac{3N^{\#} k T_o}{N \sum_{i=1} \underline{p}_i^2 / m} - 1 \right) \underline{p}_i dt \quad (15)$$

The constant γ is a coupling constant to an imaginary heat bath. When $\gamma dt = 0$ the system is effectively decoupled and Newtonian dynamics are recovered. When $\gamma dt = 1$ we obtain a constraint equation very closely related to

the momentum scaling procedure

$$d\mathbf{p}_i = \mathbf{F}_i dt + (f_o^2(t) - 1) \mathbf{p}_i \quad (16)$$

All the previous constraint equations were derived without recourse to the total Hamiltonian equations of motion for both the system and heat bath. By fixing the constraint of constant T equation (9) a truly canonical distribution of both coordinates and momenta cannot be achieved since in a truly canonical (N,V,T) ensemble the total kinetic energy is distributed according to the Boltzmann probability

$$\frac{d \sum_{i=1}^N \mathbf{p}_i^2/m}{dn} \propto \exp \left(\frac{\sum_{i=1}^N \mathbf{p}_i^2/m}{2k T_o} \right) \quad (17)$$

where n is the probability density.

This problem of simulating canonical kinetic energy fluctuations with reversible deterministic equations of motion has been solved by Nosé [6,13]. By writing down the extended system (i.e. system + heat exchanger) Hamiltonian and introducing a new scaling variable in time, Nosé was able to derive the deterministic constraint equations of motion that are both time-reversible and generate the canonical ensemble up to factors of the order D/6N where D is the number of degrees-of-freedom. The whole system may accelerate or decelerate to permit fluctuations in T without disturbing the coordinate distribution.

The Nosé ES (extended system) Hamiltonian for the system plus heat exchanger is

$$H_N = \Phi(\underline{r}_1 \dots \underline{r}_N) + \sum_i \frac{p_i^2/2m}{s^2} + (N^\# + 1) kT \log_e s + \frac{p_s^2}{2Q} \quad (18)$$

s is an arbitrary dimensionless time-scaling variable, p_s its conjugate momentum. Q (effective "mass") has units of energy time², p_s units of energy time.

The equations of motion derived by Nosé which correspond to his extended system Hamiltonian are

$$\frac{dr_i}{dt} = \frac{p_i}{ms^2} \quad (19)$$

$$\frac{dp_i}{dt} = F_i \quad (20)$$

$$d_s = \frac{p_s}{Q} dt \quad \sum_{i=1}^N \frac{p_i^2 (p_i + F_i dt)^2}{ms^3} - (N^\# + 1)kT/s \quad dt \quad (21)$$

the additional degree of freedom ($N^\# + 1$) is due to the conjugate momentum of the heat exchanger.

Hoover [12] has observed that Nosé's equations can be written in a much more amenable form by the simple expedient of reducing the global time-scale by the factor s , putting $dt = s dt'$

$$\frac{dr_i}{dt'} = \frac{p_i}{ms} \quad (22)$$

$$\frac{dp_i}{dt'} = F_i s dt' \quad (23)$$

$$ds = \frac{sp_s dt'}{Q} \quad (24)$$

$$dp_s = \left(\sum_{i=1}^N \frac{p_i^2}{ms^2} - (N^\# + 1)kT \right) dt' \quad (25)$$

re-writing the evolution equations (differentiate equation(22) with respect to time again and substitute from (23) and (24) for dp_i and ds .

The variable s is then eliminated by re-writing the second-order equation

$$\frac{d^2 r_i}{dt'^2} = \frac{1}{s} \frac{dp_i/m}{dt'} - \frac{p_i}{m_s} \frac{1}{s} \frac{ds}{dt'} \quad (26)$$

substitute from equation (24) for ds/s

$$= \frac{F_i}{m} - \frac{dr_i}{dt} \frac{p_s}{Q} \quad (27)$$

The Nosé-Hoover canonical equations of motion are then (redefining dt)

$$dr_i = p_i/m dt \quad (28)$$

$$dp_i = \frac{F_i}{m} dt - \frac{p_s}{Q} p_i dt \quad (29)$$

$$dp_s = \left(\frac{\sum_{i=1}^N p_i^2/m - NkT}{Q} \right) dt \quad (30)$$

Any choice of the constant Q will yield the canonical equations so for practical purposes it is convenient to eliminate Q also by setting

$$Q = 3NkT(\Delta t)^2 \quad (31)$$

In this case the more specific equations of motion can then be simplified into a more recognisable form

$$\underline{dp}_i = \underline{F}_i dt - \left(\frac{p_s}{3NkT\Delta t} \right) \underline{p}_i \quad (32)$$

denoting the dimensionless conjugate heat exchanger momentum as π_s then

$$d\pi_s = (f_0^2 - 1) \pi_s \quad (33)$$

is independent of the time increment in the integration of the equations of motion.

Thus the Nosé-Hoover equation in a form suitable for conversion to a finite difference algorithm is

$$d\underline{p}_i = \underline{F}_i(t) dt - \int_0^t (f_0^2 - 1) \frac{dt}{t} \underline{p}_i \quad (34)$$

Alternatively, for most practical applications imaginable, we have a practicable simple generalisation of the equations of motion

$$d\underline{p}_i = \underline{F}_i dt + (f^n - 1) \underline{p}_i \quad (35)$$

n = 0 Newtonian

n = 1 Continuous renormalisation isokinetic

n = 2 Isothermal

Example: Verlet algorithm

$$\left. \begin{aligned} m\Delta\underline{r}_i &= \underline{p}_i(t^-)\Delta t + \underline{F}_i(t)\Delta t^2 - (\pi_s + f_0^2 - 1)\underline{p}_i(t^-)\Delta t \\ \Delta\underline{p}_i &= \underline{F}_i(t)\Delta t - (\pi_s + f_0^2 - 1)\underline{p}_i(t^-) \\ \Delta\pi_s &= (f_0^2 - 1) \end{aligned} \right\} \text{Nosé-Hoover} \quad (36)$$

or for the simple generalised form

$$m\Delta\underline{r}_i = \underline{p}_i(t^-)\Delta t + \underline{F}_i(t)\Delta t^2 - (f^n - 1)\underline{p}_i(t^-)\Delta t \quad (37)$$

If anyone were to find any differences between the results obtained from amongst the above methods which significantly affect applications, e.g. comparisons with experiment or theory, the topic would be of more interest. Differences between the ensembles become negligible anyhow for a system of sufficient size.

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VECTORISING YOUR MD PROGRAM FOR THE CYBER 205
THE BRODE AND AHLRICHS WAY

by David Brown, Chemistry Department, UMIST.

As a regular reader of this illustrious, World renowned informal newsletter I have been surprised that little, if anything, has appeared in its pages concerning the adaptation of MD programs to take advantage of the "vector" capability of the CYBER 205. In Britain access to the CYBER 205 at the University of Manchester Regional Computing Centre (UMRCC) has now been possible for over two years and despite some early teething problems, with compiler bugs and fragile links with the front-end, is now providing a fairly robust service. It seems reasonable to assume then that many people have already managed to vectorise their programs but have not taken up the unique opportunity that this newsletter provides of communicating their experience directly to those most likely to be interested. I have, therefore, taken it upon myself, at least partially, and *under absolutely no duress at all from the Editor-in-chief (Bill Smith)*, to redress this balance. Hopefully others will find faults with what I'm about to relate and will be spurred on to point these out in an article of their own.

As mentioned in the title, the method I have used has been lifted wholesale from a paper by Brode and Ahlrichs [Computer Physics Communications, 42, 51, (1986)] who report that significant savings (93%) in CPU time can be achieved using it. There are, however, some slight technical differences between my own implementation and that originally described to which I will allude in due course. Only a brief resumé of the method will be given here. For a fuller description the reader is directed towards the original article.

METHOD

In a conventional MD simulation the vast majority of the time is spent evaluating the forces between all possible interacting pairs. Schematically the conventional order of evaluating pair separations ($r_{ij}=r_i-r_j$), and hence forces, is shown below for a system of just $N=6$ particles.

	J=1	2	3	4	5	6
I=1	r_{11}	r_{12}	r_{13}	r_{14}	r_{15}	r_{16}
2	r_{21}	r_{22}	r_{23}	r_{24}	r_{25}	r_{26}
3	r_{31}	r_{32}	r_{33}	r_{34}	r_{35}	r_{36}
4	r_{41}	r_{42}	r_{43}	r_{44}	r_{45}	r_{46}
5	r_{51}	r_{52}	r_{53}	r_{54}	r_{55}	r_{56}
6	r_{61}	r_{62}	r_{63}	r_{64}	r_{65}	r_{66}

This is achieved in Fortran quite simply within a double loop

in which the first index, I, runs from 1 to N-1 and the second, J, runs from I+1 to N. On a scalar processor this represents an easy and efficient way of calculating the matrix for a reasonable number of particles, i.e. N of order 10^2 . On the CYBER 205 though there are distinct disadvantages with this approach. Firstly, the vector lengths are short, at most N-1 and, secondly, they become shorter as I→N. The CYBER 205, however, only starts to approach its asymptotic speed when vector lengths are of order 10^3 , due to penalties incurred in starting up an operation. Now consider the following scheme where the interaction matrix and an identical copy have been juxtaposed

	J=1	2	3	4	5	6	7	8	9	10	11	12
I=1	r ₁₁	r ₁₂	r ₁₃	r ₁₄	r ₁₅	r ₁₆	r ₁₁	r ₁₂	r ₁₃	r ₁₄	r ₁₅	r ₁₆
2	r ₂₁	r ₂₂	r ₂₃	r ₂₄	r ₂₅	r ₂₆	r ₂₁	r ₂₂	r ₂₃	r ₂₄	r ₂₅	r ₂₆
3	r ₃₁	r ₃₂	r ₃₃	r ₃₄	r ₃₅	r ₃₆	r ₃₁	r ₃₂	r ₃₃	r ₃₄	r ₃₅	r ₃₆
4	r ₄₁	r ₄₂	r ₄₃	r ₄₄	r ₄₅	r ₄₆	r ₄₁	r ₄₂	r ₄₃	r ₄₄	r ₄₅	r ₄₆
5	r ₅₁	r ₅₂	r ₅₃	r ₅₄	r ₅₅	r ₅₆	r ₅₁	r ₅₂	r ₅₃	r ₅₄	r ₅₅	r ₅₆
6	r ₆₁	r ₆₂	r ₆₃	r ₆₄	r ₆₅	r ₆₆	r ₆₁	r ₆₂	r ₆₃	r ₆₄	r ₆₅	r ₆₆

This time all possible interactions have been taken, as $r_{ij} = -r_{ji}$, but the vector length has been kept to N except, in this case, for the last one which equals N/2. Of course, there are always (N-1)/2 blocks of N interactions, so for an odd number of particles there would be an integer number of blocks. This kind of procedure can be realised on the computer by doubling the length of the position arrays, such that $r_{i+N} = r_i$, and also the force arrays. A double loop could then be performed in which the outer loop controller takes values of the difference in the indices of the arrays $J-I=1,2,3,\dots$ etc. whilst the inner one runs from 1 to N or until all possible interactions have been taken. Then the total force will be given by $f_i + f_{i+N}$. In practice several blocks of N interactions are processed at the same time using additional vectors of length ~5000 so as to improve efficiency.

FORTRAN 200 CODE

To demonstrate the simplicity of this method selected sections of a "typical" MD code for an equilibrium simulation of a single component atomic system, interacting through a Lennard-Jones 12-6 potential, will be given.

The positions of the N particles reside in the arrays X, Y and Z of dimension N2=2*N and are contained in a periodic cube such that all coordinates lie within the range -FAC and +FAC. The forces will be accumulated in the vectors FX, FY and FZ again of length N2. All the other vectors used V1, V2, V3, V4, R2, XIJ, YIJ, ZIJ including the BIT VECTOR BV are of length NVECL where NVECL is usually, though not necessarily, a multiple of N in the region of 5000. Use of FORTRAN 200 shorthand has been used throughout so statements such as

```
A(1;N)=B(1+N;N)**2
```

are short for

```
DO 10 I=1,N  
  A(I)=B(I+N)**2  
10 CONTINUE
```

Another way this can be written makes use of array DESCRIPTORS which also have to be declared. These are all suffixed with a "D" so as to avoid confusion with scalars. A descriptor contains a starting address and a length both set by an ASSIGN statement. The example above would then look like

```
ASSIGN AD,A(1;N)  
ASSIGN BD,B(1+N;N)  
AD=BD**2
```

WARNING!! As the code given below has been lifted and adapted from more complicated programs it will not in itself have been tested on the computer. There remains, therefore, the possibility, although very slight, that it contains one or two small buggettes.

```
:  
:  
C   Declare arrays, descriptors and bit vectors etc.  
C  
  REAL X(N2),Y(N2),Z(N2),FX(N2),FY(N2),FZ(N2)  
  REAL XIJ(NVECL),YIJ(NVECL),ZIJ(NVECL),R2(NVECL),  
+V1(NVECL),V2(NVECL),V3(NVECL),V4(NVECL)  
  BIT BV(NVECL),BVD  
  DESCRIPTOR BVD,XIJD,YIJD,ZIJD,R2D,V1D,V2D,V3D,V4D  
:  
:  
:  
C  
C   Calculate from the number of interactions, NITCC,  
C   the maximum number of blocks of N that can be processed  
C   at any one time, NMAXCC, and the number of batches  
C   required to account for all interactions, NTIMCC.  
C  
  NITCC=NINT(0.5*N*(N-1))  
  NMAXCC=NVECL/N  
  MAXCC=NMAXCC*N  
  NTIMCC=NITCC/MAXCC  
  MTIMCC=NTIMCC*MAXCC  
  IF(MTIMCC.NE.NITCC)NTIMCC=NTIMCC+1  
:  
:  
:  
C  
C   Zero potential energy and force accumulators.  
C  
  TCCPEN=0.0  
  FX(1;N2)=0.0  
  FY(1;N2)=0.0  
  FZ(1;N2)=0.0  
C
```

```

C      Normalize coordinates using RFAC=1.0/FAC so that they
C      lie between -1 and +1.
C
C      X(1;N)=X(1;N)*RFAC
C      Y(1;N)=Y(1;N)*RFAC
C      Z(1;N)=Z(1;N)*RFAC
C
C      Copy coordinates onto the 1+N to N2 elements.
C
C      X(1+N;N)=X(1;N)
C      Y(1+N;N)=Y(1;N)
C      Z(1+N;N)=Z(1;N)
C
C      Zero counters and start loop over the number of batches.
C
C      L=0
C      DO 2000 K=1,NTIMCC
C      LL=0
C
C      Calculate the present value of the difference in
C      indices, JI=J-I.
C
C      JI=(K-1)*NMAXCC
C
C      Loop over the number of blocks of N interactions in one
C      batch
C
C      DO 2010 M=1,NMAXCC
C
C      Increment index difference counter
C
C      JI=JI+1
C
C      Calculate the number of interactions in this set, JMAX,
C      allowing for the possibility that JMAX=N/2 or 0.
C
C      JMAX=NITCC-L
C      JMAX=MINO(N,JMAX)
C      IF(JMAX.EQ.0) GOTO 2030
C
C      Store separation vectors in elements offset by LL
C
C      XIJ(LL+1;JMAX)=X(1;JMAX)-X(JI+1;JMAX)
C      YIJ(LL+1;JMAX)=Y(1;JMAX)-Y(JI+1;JMAX)
C      ZIJ(LL+1;JMAX)=Z(1;JMAX)-Z(JI+1;JMAX)
C
C      Update offsets
C
C      L=L+JMAX
C      LL=LL+JMAX
C
C      2010 CONTINUE
C      2030 CONTINUE
C
C      Processing the LL interactions stored
C
C      ASSIGN BVD,BV(1;LL)
C      ASSIGN R2D,R2(1;LL)
C      ASSIGN XIJD,XIJ(1;LL)
C      ASSIGN YIJD,YIJ(1;LL)

```

```

ASSIGN ZIJD,ZIJ(1;LL)
ASSIGN V4D,V4(1;LL)
C
C Apply periodic boundaries making use of vector function
C VAINT which returns into the target vector, in this
C case V4D, the integer part of the subject vector as a
C real.
C
XIJD=XIJD-2.0*VAINT(XIJD;V4D)
YIJD=YIJD-2.0*VAINT(YIJD;V4D)
ZIJD=ZIJD-2.0*VAINT(ZIJD;V4D)
C
C Calculate the square of the separation
C
R2D=XIJD**2+YIJD**2+ZIJD**2
C
C Set the elements of the bit vector to 1 if r2 is
C less than the square of the potential cutoff, CUTSQ.
C
BVD=R2D.LT.CUTSQ
C
C Count the number of interactions within cutoff using
C special Q8 routine, Q8SCNT.
C
NN=Q8SCNT(BVD)
C
C Check for no contributing interactions
C
IF(NN.EQ.0) GOTO 2000
C
C Assign some descriptors to new length NN
C
ASSIGN V1D,V1(1;NN)
ASSIGN V2D,V2(1;NN)
ASSIGN V3D,V3(1;NN)
ASSIGN V4D,V4(1;NN)
C
C Use the Q8VCMFRS instruction to compress the list of
C square separations down to just those that contribute.
C
V2D=Q8VCMFRS(R2D,BVD;V2D)
C
C Re-assign array descriptor.
C
ASSIGN R2D,R2(1;NN)
C
C Calculate terms in potential and force using
C  $CON1=\sigma^2/FAC**2$  and  $CON2=48*FAC*\epsilon/\sigma^2$ 
C
R2D=CON1/V2D
V2D=R2D**3
V1D=V2D**2
V3D=V1D-V2D
V4D=CON2*R2D*(V1D-0.5*V2D)
C
C Accumulate total potential energy (/4 $\epsilon$ ) using the Q8SSUM
C instruction.
C
TCCPEN=TCCPEN+Q8SSUM(V3D)

```

```

C
C   Compress down the components of the separation vector
C
V1D=Q8VCMPRS(XIJD,BVD;V1D)
V2D=Q8VCMPRS(YIJD,BVD;V2D)
V3D=Q8VCMPRS(ZIJD,BVD;V3D)
C
C   Re-assign separation descriptors
C
ASSIGN XIJD,XIJ(1;NN)
ASSIGN YIJD,YIJ(1;NN)
ASSIGN ZIJD,ZIJ(1;NN)
C
C   Calculate components of the force
C
XIJD=V1D*V4D
YIJD=V2D*V4D
ZIJD=V3D*V4D
C
C   Accumulate the potential contribution to the pressure
C   tensor using the Q8SDOT instruction which returns the
C   sum of the dot product of two vectors.
C
A11=A11+Q8SDOT(V1D,XIJD)
A12=A12+Q8SDOT(V1D,YIJD)
A13=A13+Q8SDOT(V1D,ZIJD)
A22=A22+Q8SDOT(V2D,YIJD)
A23=A23+Q8SDOT(V2D,ZIJD)
A33=A33+Q8SDOT(V3D,ZIJD)
C
C   Re-assign descriptors for expansion.
C
ASSIGN V1D,V1(1;LL)
ASSIGN V2D,V2(1;LL)
ASSIGN V3D,V3(1;LL)
C
C   Using the Q8VXPND instruction the force arrays are
C   expanded up to their original length using the control
C   bit vector BV to insert zeroes into non-contributing
C   interactions.
C
V1D=Q8VXPND(XIJD,BVD;V1D)
V2D=Q8VXPND(YIJD,BVD;V2D)
V3D=Q8VXPND(ZIJD,BVD;V3D)
C
C   Calculate number of interactions already processed.
C
LOLD=L-LL
C
C   Assign forces to particles using the reverse procedure
C   as that used above to assign separations and Newton's
C   law  $f_{ij}=-f_{ji}$ .
C
LL=0
JI=(K-1)*NMAXCC
DO 2120 M=1,NMAXCC
JI=JI+1
JMAX=NITCC-LOLD
JMAX=MINO(N,JMAX)

```

```

IF(JMAX.EQ.0) GOTO 2130
FX(1;JMAX)=FX(1;JMAX)+V1(LL+1;JMAX)
FY(1;JMAX)=FY(1;JMAX)+V2(LL+1;JMAX)
FZ(1;JMAX)=FZ(1;JMAX)+V3(LL+1;JMAX)
FX(JI+1;JMAX)=FX(JI+1;JMAX)-V1(LL+1;JMAX)
FY(JI+1;JMAX)=FY(JI+1;JMAX)-V2(LL+1;JMAX)
FZ(JI+1;JMAX)=FZ(JI+1;JMAX)-V3(LL+1;JMAX)
LOLD=LOLD+JMAX
LL=LL+JMAX
2120 CONTINUE
2130 CONTINUE
2000 CONTINUE
C
C   End of force calculation.
C
C   Add forces in array elements I and I+N to obtain
C   total force on particle I.
C
FX(1;N)=FX(1;N)+FX(1+N;N)
FY(1;N)=FY(1;N)+FY(1+N;N)
FZ(1;N)=FZ(1;N)+FZ(1+N;N)
C
:
:

```

It should be noted that with the vectorisation of the forces loop the rest of the integration algorithm can start to account for a sizeable percentage of the overall CPU time used. It is, therefore, essential that all these other parts of the code are also vectorised. This is generally straightforward as they usually consist of just N-particle loops.

DISCUSSION

The scheme described differs slightly from that of the originators in three main respects :-

1) The separation vectors are compressed before multiplying by V4 to obtain the resultant forces instead of expanding V4 and then multiplying by XIJ etc.

Although more compressions and expansions are required this process represents a significant saving in CPU time, particularly if the pressure tensor is being calculated.

2) Exact force and potential calculations are carried out rather than using tabulated forces and potentials.

For inverse even power potentials there is little saving using tables by the time all the gathering and interpolating has been done. If constant temperature algorithms are being used in conjunction with constant shear rate or stress methods a high degree of accuracy is required in the force calculation to prevent a drift in the constrained properties.

3) Calculation of $g(r)$ is omitted from the main loop.

Accumulating the radial distribution function presents quite a problem as it is inherently non-vectorisable due to the random access of array elements. It needs be a separate

loop could be constructed containing the necessary Fortran and accessed periodically instead of the above loop. As configurations tend to be well correlated for many steps there is little point accumulating $g(r)$ too often. For my own part I prefer to store coordinates etc. on magnetic tape for processing later and calculate $g(r)$ if and when necessary.

Another point to note is that correlation functions have also been calculated from data stored on magnetic tape using the Fourier transform method described by Smith [see CCP5 Newsletters 5 and 7]. This is a very efficient way of calculating correlation functions as the fast Fourier transform (FFT) library of routines on the CYBER 205 is also vectorized.

In all, the method described results in increases of speed by a factor of about 15 compared to a scalar optimized program run on either the CDC 7600 or CYBER 176 machines also at UMRCC (though, sadly, soon to depart for that big computer centre in the sky). Although this comparison does not include the extra time needed to calculate $g(r)$. Even so the increase in speed is at least an order of magnitude, which is not to be sneezed at. It is also possible that the above scheme may improve performance figures on CRAY machines as well. This may have already been tried by someone, in which case I would like to hear from them.

MINIDAP AT DAPSU

J.E.Quinn
DAPSU
Queen Mary College

DAPSU (DAP Support Unit) took delivery of a prototype Mini-DAP as part of a collaborative grant between SERC and ICL during April of last year. Since then the mini-DAP has been in daily use, producing demonstrations for a variety of projects including hydrodynamics, molecular graphics, finite element work and for general development.

DAPs (Distributed Array Processors) are already well known in the molecular simulation field and various articles describing their use have appeared in this newsletter and other publications (see for example references at end). Essentially a DAP is a collection of individual processors all of which are simultaneously able to execute the same instruction. Each processor has its own memory but also has links with its four nearest neighbours and 'highways' are present for broadcasting data simultaneously to all processors. In addition each processor has its own on/off switch, which is used, for example, to map directly onto conditional clauses in user programs, thus implementing them very efficiently.

Mini-DAP

Mini-DAP grew out of a previous version of the DAP, one of which was installed at QMC in 1980 and has been used by the scientific community worldwide, though mainly in this country, since then. The term mini-DAP is deserved for two reasons;

- 1) The original DAP at QMC was built in MSI (medium scale integration) technology and occupied several double filing cabinet size boxes, whereas the mini-DAP is built in LSI (large scale integration) and fits under a table.
- 2) MSI DAPs have 4096 processors, mini-DAP has 1024.

The new production version of DAP is primarily intended for use in a single user workstation environment with the DAP as an attached processor to either a SUN or a microVAX. The prototype at QMC is currently attached to a PERQ II.

Although the DAP was originally manufactured by ICL, since November 1986 the DAP has been placed in the control of a new company called 'Active Memory Technology' (AMT) who are exploiting the DAP commercially. AMT has close links with DAPSU (several ex DAPSU members have become founder members of AMT) and many of the current mini-DAP demonstrations to be seen at computer exhibitions were written by DAPSU staff.

Mini-DAP Software Environment

Mini-DAP has a very similar software environment to the MSI DAP, although there are four significant changes,

- 1) The array size is 1024 instead of 4096
- 2) The data area in DAP and HOST are no longer shared
- 3) A fast I/O channel (FIO) operating at rates up to 40 Mbytes/sec is available
- 4) Asynchronous working of DAP and HOST is possible.

1) This change is fundamental and will require many of the existing DAP programs which make specific reference to either the total number of processors (4096) or edge size (64) of the MSI DAP, to be altered. The syntax of the DAP FORTRAN declarations for vector and matrix working, e.g.

```
REAL*3 R( ), S( )
```

remains the same, but the value of the constrained dimensions is 32 for the mini-DAP.

2) Whereas the MSI DAP used common store for DAP and HOST sections of the program, mini-DAP requires data to be sent to, or received from the DAP. This transfer is actually performed by the HOST program and two routines DAPSEN and DAPREC are provided for this purpose. Conversion of data to DAP storage mode must still be carried out as on the MSI DAP.

3) The fast I/O provides new ways of accessing the results of a DAP run, or of inputting data quickly. For example, a video board has been designed to sit on the FIO enabling the DAP store to be used in the same way as a conventional frame store, but at a much higher rate. The DAP FIO is asynchronous with the working of the DAP, so that data can be fed in or read out whilst DAP processing continues.

4) The ability to use DAP and HOST machines asynchronously creates several programming possibilities. Some of the current demonstrations which use external graphics processors to display their results have the DAP calculating the next picture, while the HOST is displaying the current one. Semaphore systems have been devised to allow either machine to poll the other to find if it has finished the next step in the calculations.

At DAPSU the prototype mini-DAP has a full set of system software allowing programs to be compiled, ran and interactively 'debugged'. In addition DAPSU staff have been converting routines from the MSI DAP subroutine library for use on the mini-DAP and ~100 routines will be available in the first release.

ICL, who have rights to mini-DAP for defence contracts in this country are supplying software for use with the video board and Intercept Systems, a software house are supplying software for a variety of DAP related problems. Further graphics work will continue at DAPSU.

Performance

It is generally reckoned that the MSI DAP was approximately one third the power of the CRAY 1S for computationally intensive calculations. Our own estimate in DAPSU show the prototype mini-DAP to be one third the power of the MSI DAP and based on discussions with AMT we believe the production machines will be one half the power of the MSI DAP, 1/6th of the CRAY 1S. Memory sizes on production machines will range from 4-128 Mbyte.

Support for mini-DAP

At present there is no front line support for mini-DAPs. However anyone wishing to use the mini-DAP in a 'pump-priming' fashion should write to Carol Hirst at DAPSU explaining what they would like to do, so that, where possible, time and resources can be made available.

The possibility of installing a large (32 Mbyte) mini-DAP at QMC and using it to provide a national service, replacing the current ICL 2980/DAP service, is being explored. At this point in time the choice of HOST is unclear as is a start date for such a service. However it is envisaged that the two machines would run in parallel for a number of months before the 2980/DAP is removed, allowing users time to move from one DAP to the other.

In addition to this service a number of research groups are applying to various grant awarding bodies for the provision of their own mini-DAP.

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See for example

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CCP5 LITERATURE SURVEY 1986.

We present in the following pages the CCP5 literature survey for the year 1986. The survey was conducted with the help of the INSPEC service at the University of London Computing Centre and we thank Dr. J. Altmann of U.L.C.C. for her essential contribution. Thanks are also due to Mrs. Anne Ireland of the Daresbury Reprographics Division for typing and sorting the entries appearing here and to Mrs. C.M. Smith for helping with the proof reading.

We invite readers who have knowledge of other relevant papers, articles or books to send us the details, which we shall be pleased to publish in an Addendum in our next newsletter.

W. Smith

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J. C. C. SHARP
DARESBURY LABORATORY